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Projectile fragmentation as a tool  
to observe shell effects close to  
the neutron drip-line

EXON 2014, Kaliningrad  
EXON 2014, Kaliningrad

09 / 8-13 / 2012

LISE++

LISE++

## 1. Introduction

- ◆ EXON2009
- ◆ EXON2012



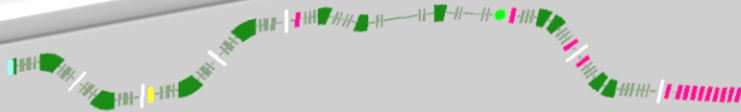
## 2. "Calcium anomaly"

## 3. Abrasion-Ablation

## 4. More probable prefragments

## 5. dBE systematic

## 6. Summary



## ☐ Search for new isotopes

- ✓ The limits of nuclear stability provide a key benchmark of nuclear models
- ✓ The context of astrophysics
  - *Understanding the r-process abundance patterns of elements*

## ☐ Production mechanism

- ✓ Production cross sections, Momentum distributions, Reaction choice
- ✓ Secondary beam intensities. Planning new experiments, set-ups (FRIB, RIBF, FAIR)

## ☐ Nuclear structure

- ✓ Changes in the structure of neutron-rich nuclei

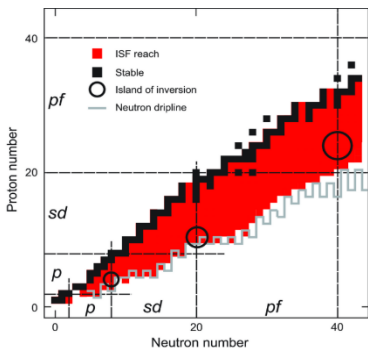
- *Region around  $^{31}\text{Na}$  is now known as the “island of inversion”*
- *Deformation around neutron number  $N = 40$  in Fe and Cr nuclei*

- ✓ Shell closures at  $N=32$  and  $N=34$

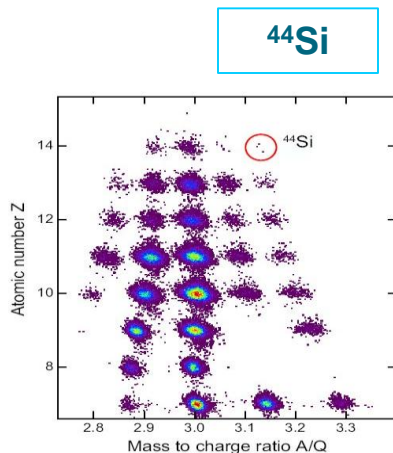
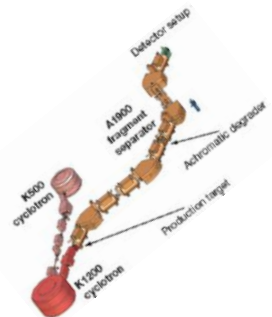
- *Masses of exotic calcium isotopes pin down nuclear forces ( $^{53,54}\text{Ca}$ )*  
F. Wienholtz et al., Nature 498 (2013) 346
- *Evidence for a new nuclear ‘magic number’ from the level structure of  $^{54}\text{Ca}$*   
D. Steppenbeck et al., Nature 502 (2013) 207

These measurements of the  $\text{Ex}(2^+_1)$  in  $^{54}\text{Ca}$  at RIKEN found that the experimental value is 0.5 MeV smaller than the prediction of the full  $pf$  shell-model space with the GXPF1B effective interaction.

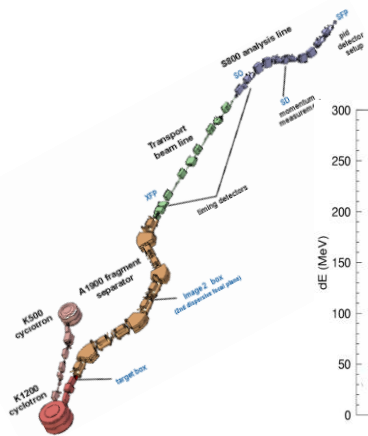
GXPF1B: Y. Utsuno, T. Otsuka, B. A. Brown, M. Honma, T. Mizusaki, and N. Shimizu,  
PRC 86, 051301(R) (2012) ; GXPF1B5 – modified GXPF1B for 0.5 MeV shift



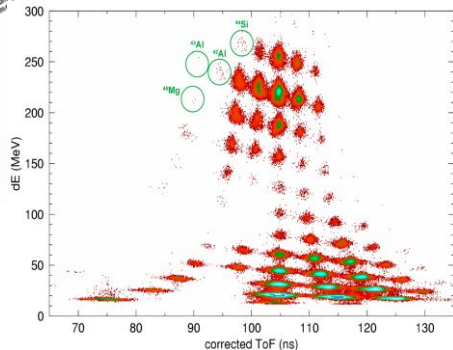
O.T. et al., Phys.Rev. C 75, 064613 (2007)



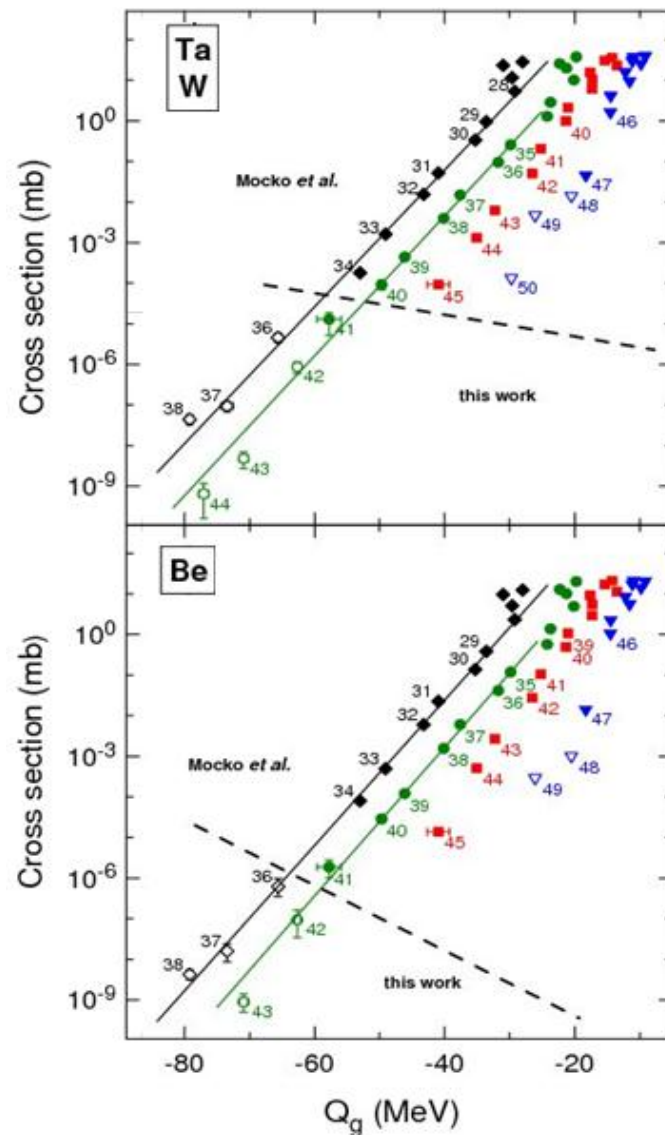
T.Baumann et al., Nature (London) 449, 1022 (2007)



**$^{40}\text{Mg}$ ,  $^{42}\text{Al}$ ,  $^{43}\text{Al}$**



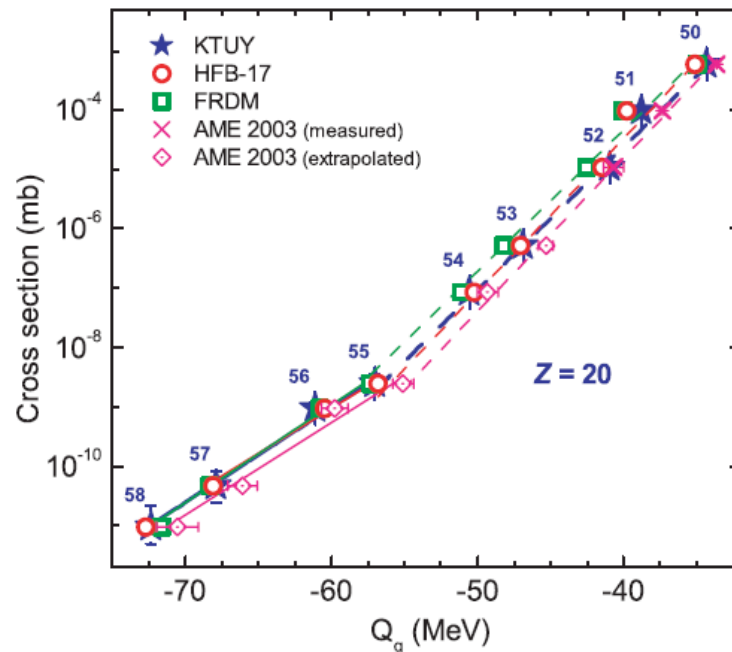
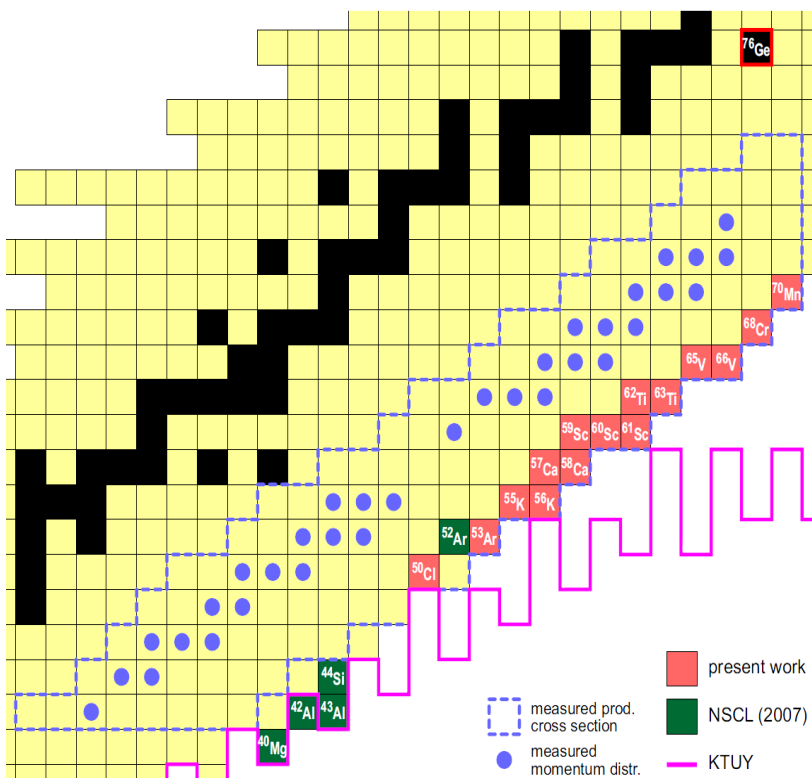
$$Q_g = ME(Z = 20, A = 48) - ME(Z, A)$$



**$^{48}\text{Ca}(140\text{MeV/u}) + \text{W, Be}$**

O.T. et al., Phys.Rev. C 75, 064613 (2007)

$^{50}\text{Cl}$ ,  $^{53}\text{Ar}$ ,  $^{55,56}\text{K}$ ,  $^{57,58}\text{Ca}$ ,  $^{59,60,61}\text{Sc}$ ,  
 $^{62,63}\text{Ti}$ ,  $^{65,66}\text{V}$ ,  $^{68}\text{Cr}$ ,  $^{70}\text{Mn}$



Enhanced cross sections might be the result of increased binding

This region (around  $^{62}\text{Ti}$ ) was previously predicted to be a new island of inversion  
*B. A. Brown Prog. Part. Nucl. Phys. 47 (2001) 517*

Phys.Rev.Lett. 102, 142501 (2009) :  
 Phys.Rev.C. 80, 034609 (2009) :  
 NIM A 620, 578-584 (2010) :

New isotopes, Evidence for a Change in the Nuclear Mass Surface  
 Set-up, cross sections, momentum distributions  
 A new approach to measure momentum distributions

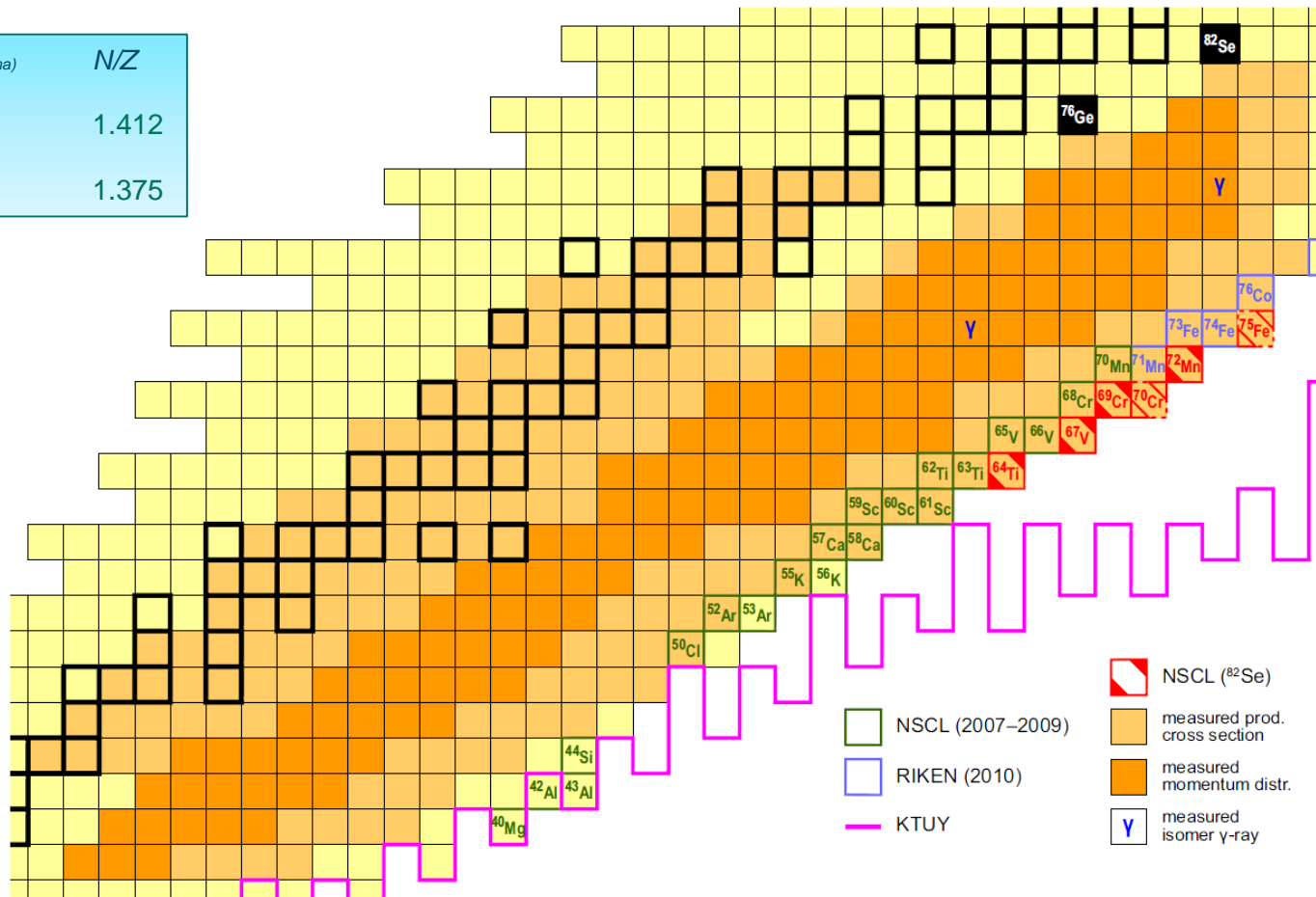
So, what is "Calcium anomaly" ?  
 Reaction property or  
 Nuclear structure feature?  
 It should be checked with another beam

$^{64}\text{Ti}$ ,  $^{67}\text{V}$ ,  $^{69}\text{Cr}$ ,  $^{72}\text{Mn}$   
 $^{70}\text{Cr}$  1event &  $^{75}\text{Fe}$  1event

$N(\sigma_W) = 90$   
 $N(\sigma_{Be}) = 330$   
 $N(d\sigma/dp) = 126$

Beam	$E$ (MeV/u)	$I$ (pna)	$N/Z$
$^{82}\text{Se}$	139	35	1.412
$^{76}\text{Ge}$	130	20	1.375

$\Delta N / \Delta Z = 2$



- NSCL ( $^{82}\text{Se}$ )
- NSCL (2007–2009)
- RIKEN (2010)
- KTUY
- measured prod. cross section
- measured momentum distr.
- Y measured isomer  $\gamma$ -ray

## Confirmation of the Calcium anomaly.... Without explanation

## 2013 : Calcium anomaly as shell effects close to drip line

 Selected for a [Viewpoint](#) in *Physics*

PHYSICAL REVIEW C **87**, 054612 (2013)



### Production cross sections from $^{82}\text{Se}$ fragmentation as indications of shell effects in neutron-rich isotopes close to the drip-line

O. B. T.,<sup>1,\*</sup> M. Portillo,<sup>2</sup> D. J. Morrissey,<sup>1,3</sup> A. M. Amthor,<sup>2</sup> L. Bandura,<sup>2</sup> T. Baumann,<sup>1</sup> D. Bazin,<sup>1</sup> J. S. Berryman,<sup>1</sup> B. A. Brown,<sup>1,4</sup> G. Chubarian,<sup>5</sup> N. Fukuda,<sup>6</sup> A. Gade,<sup>1,4</sup> T. N. Ginter,<sup>1</sup> M. Hausmann,<sup>2</sup> N. Inabe,<sup>6</sup> T. Kubo,<sup>6</sup> J. Pereira,<sup>1</sup> B. M. Sherrill,<sup>1,4</sup> A. Stolz,<sup>1</sup> C. Sumithrarachichi,<sup>1</sup> M. Thoennessen,<sup>1,4</sup> and D. Weisshaar<sup>1</sup>

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<sup>2</sup>Facility for Rare Isotope Beams, Michigan State University, East Lansing, Michigan 48824, USA

<sup>3</sup>Department of Chemistry, Michigan State University, East Lansing, Michigan 48824, USA

<sup>4</sup>Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan 48824, USA

<sup>5</sup>Cyclotron Institute, Texas A&M University, College Station, Texas 77843, USA

<sup>6</sup>RIKEN Nishina Center, RIKEN, Wako-shi, Saitama 351-0198, Japan

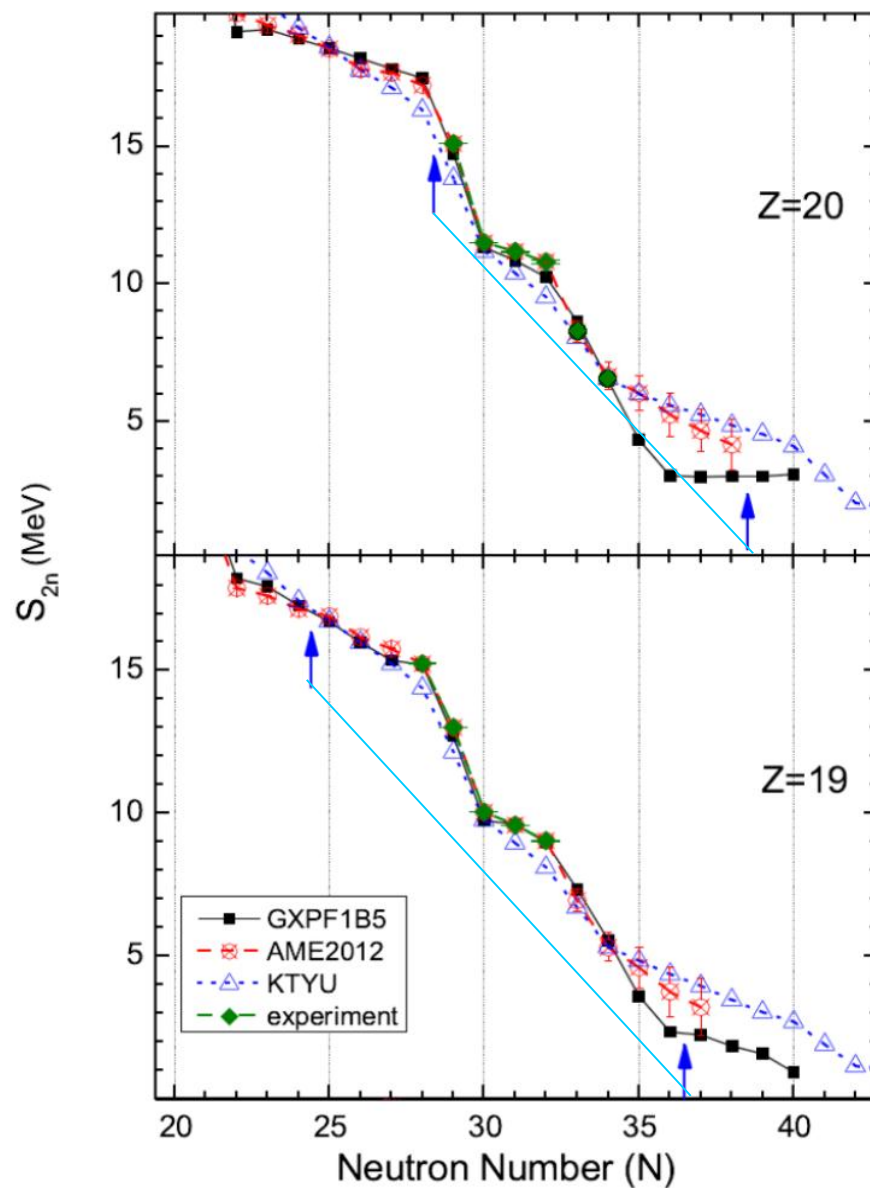
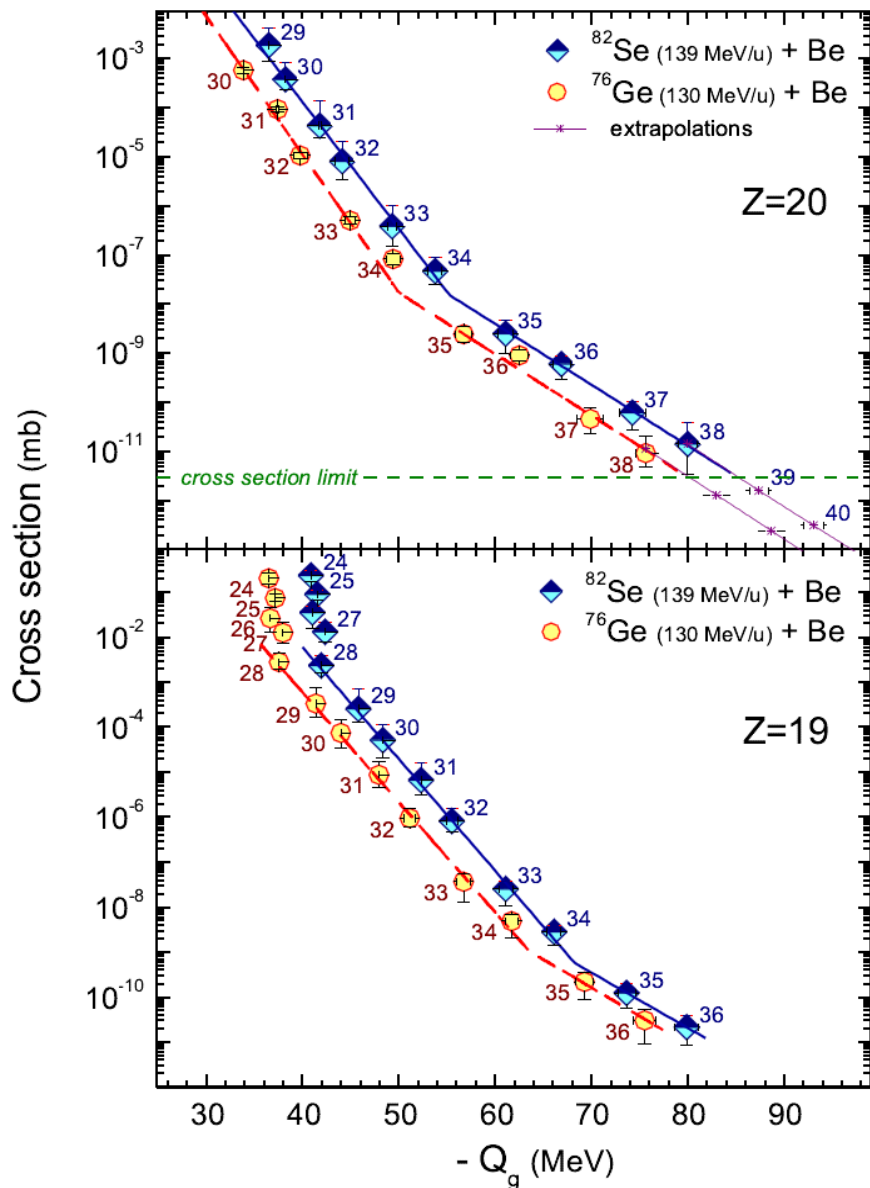
The measured cross sections were best reproduced by using masses derived from the full  $pf$  shell-model space with the GXPF1B5 [1] effective interaction modified to a recent  $^{54}\text{Ca } E_x(2^+_{-1})$  measurement [2].

The “Calcium anomaly” can be explained with a shell model that predicts a subshell closure at  $N = 34$  around  $Z = 20$ .

[1] M. Honma, T. Otsuka, B. A. Brown, and T. Mizusaki, *Eur. Phys. J. A* **25**, Suppl. 1, 499 (2005)

[2] D. Steppenbeck et al., *Nature* **502**, 207 (2013)





Experimental masses:

A.T. Gallant et al., Phys. Rev. Lett. 109, 032506 (2012)

F. Wienholtz et al., Nature 498 (2013) 346

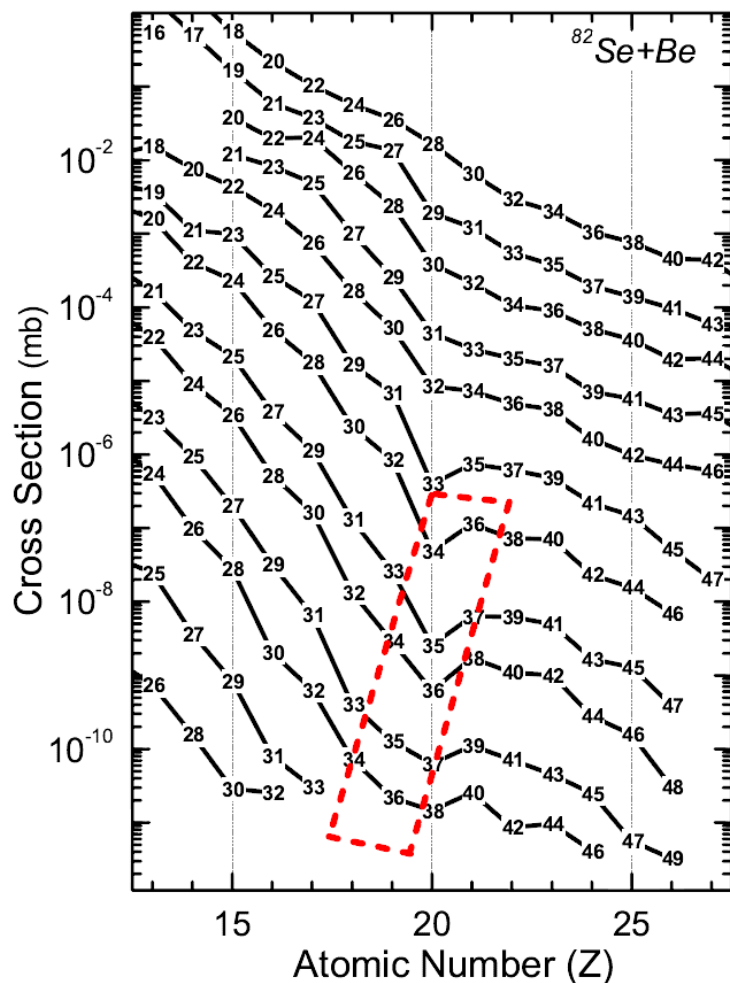


FIG. Production cross section versus atomic number ( $Z$ ) for fragments from reaction of  $^{82}\text{Se}$  with beryllium targets. Lines are connected according to constant  $N - 2Z$ , while labels represent the neutron number.

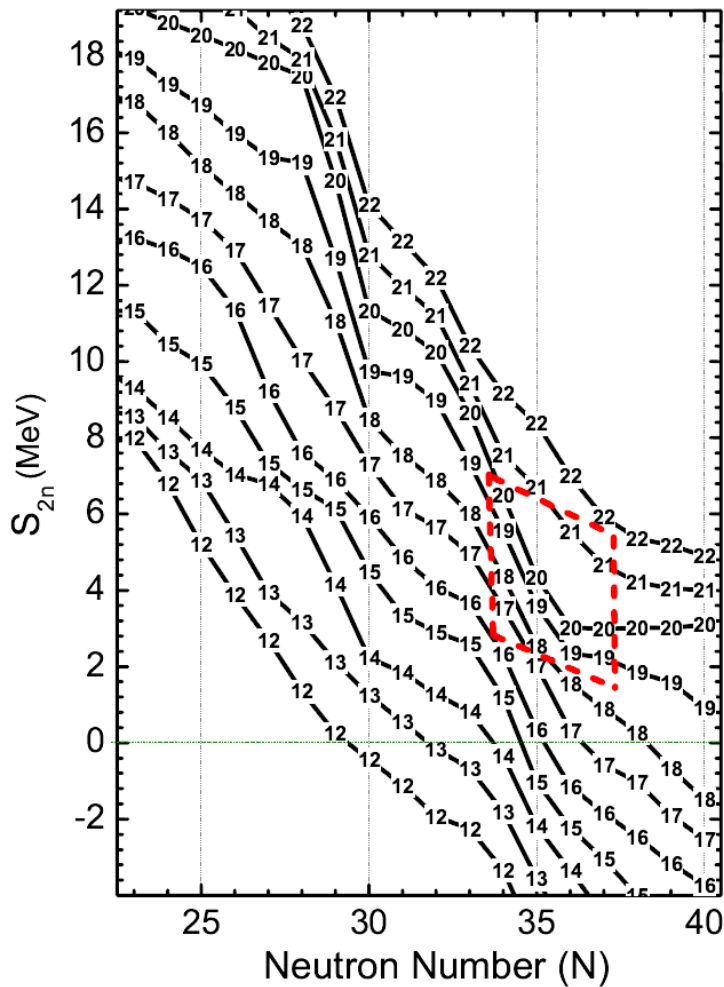


FIG. Two-neutron separation energy  $S_{2n}$  versus neutron number ( $N$ ) for elements  $12 \leq Z \leq 22$ . Values are calculated using results from the GXPF1B5 model. Labels in the lines show atomic numbers of nuclei.

No such dump with other theoretical models

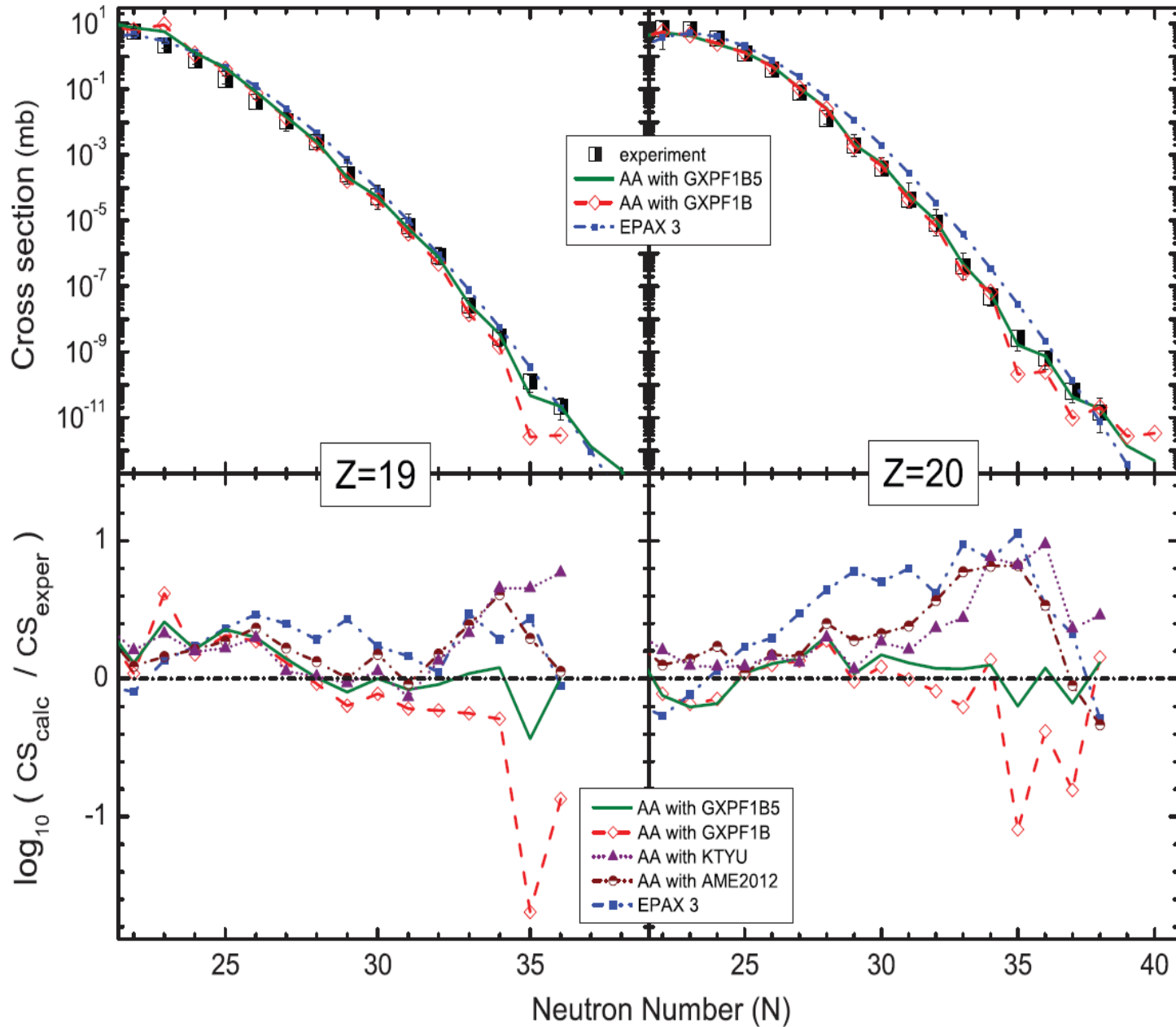


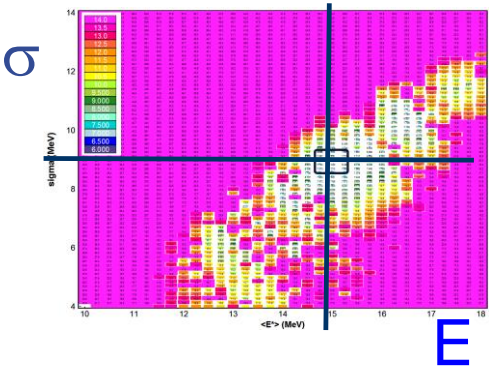
The Abrasion-Ablation model is very sensitive to the input mass values for the most exotic nuclei.

Different mass models as

- HFB9,
- HFB17,
- KTYU
- TUYU
- AME2003
- AME2012
- GXPF1B,
- GXPF1B5

were used in LISE<sup>++</sup> Abrasion-Ablation excitation energy minimization procedures to compare with the experimental data





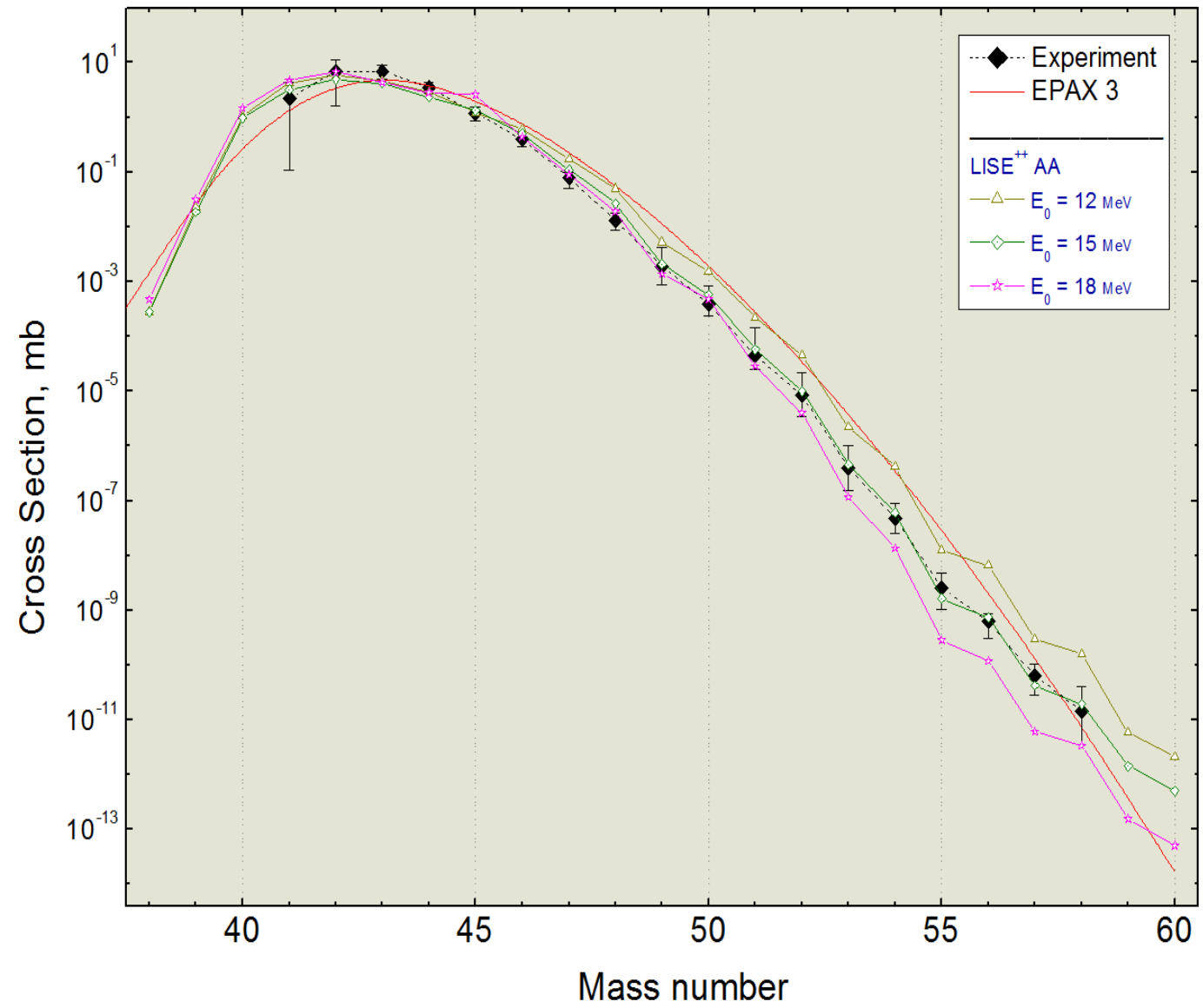
The best result to describe the experimental data of isotopes of elements  $16 < Z < 24$  has been obtained with GXPF1B5 (+ LDM0) at  $E^* = 15.0$  ( $\sigma = 9.15$ ) MeV

GSI:

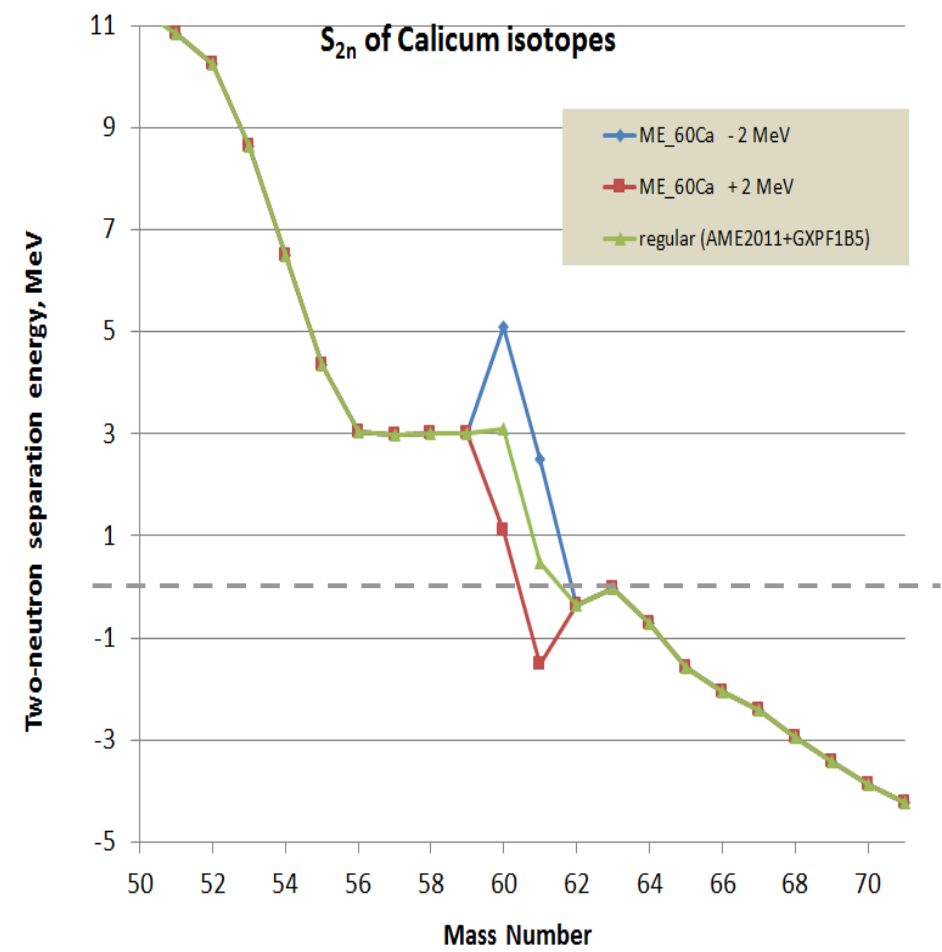
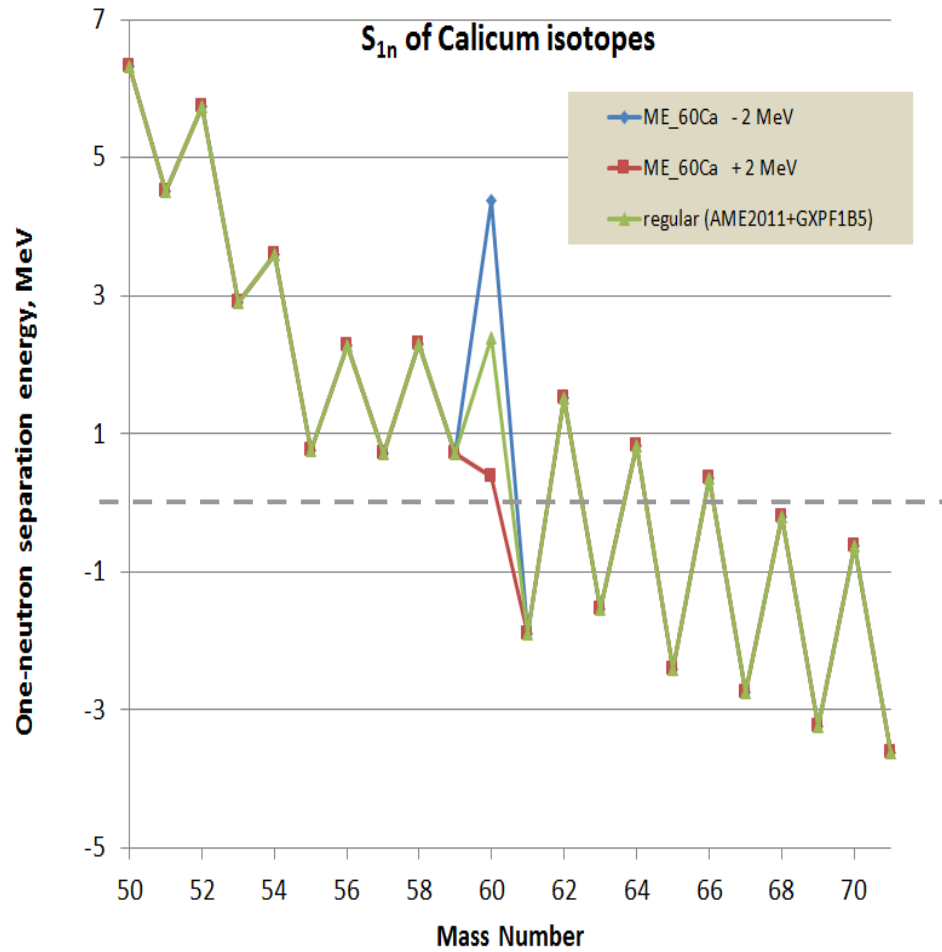
$^{40}\text{Ar}$  beam:  
 $\langle E \rangle = 13.3$  MeV  
 NPA 531,709 (1991)

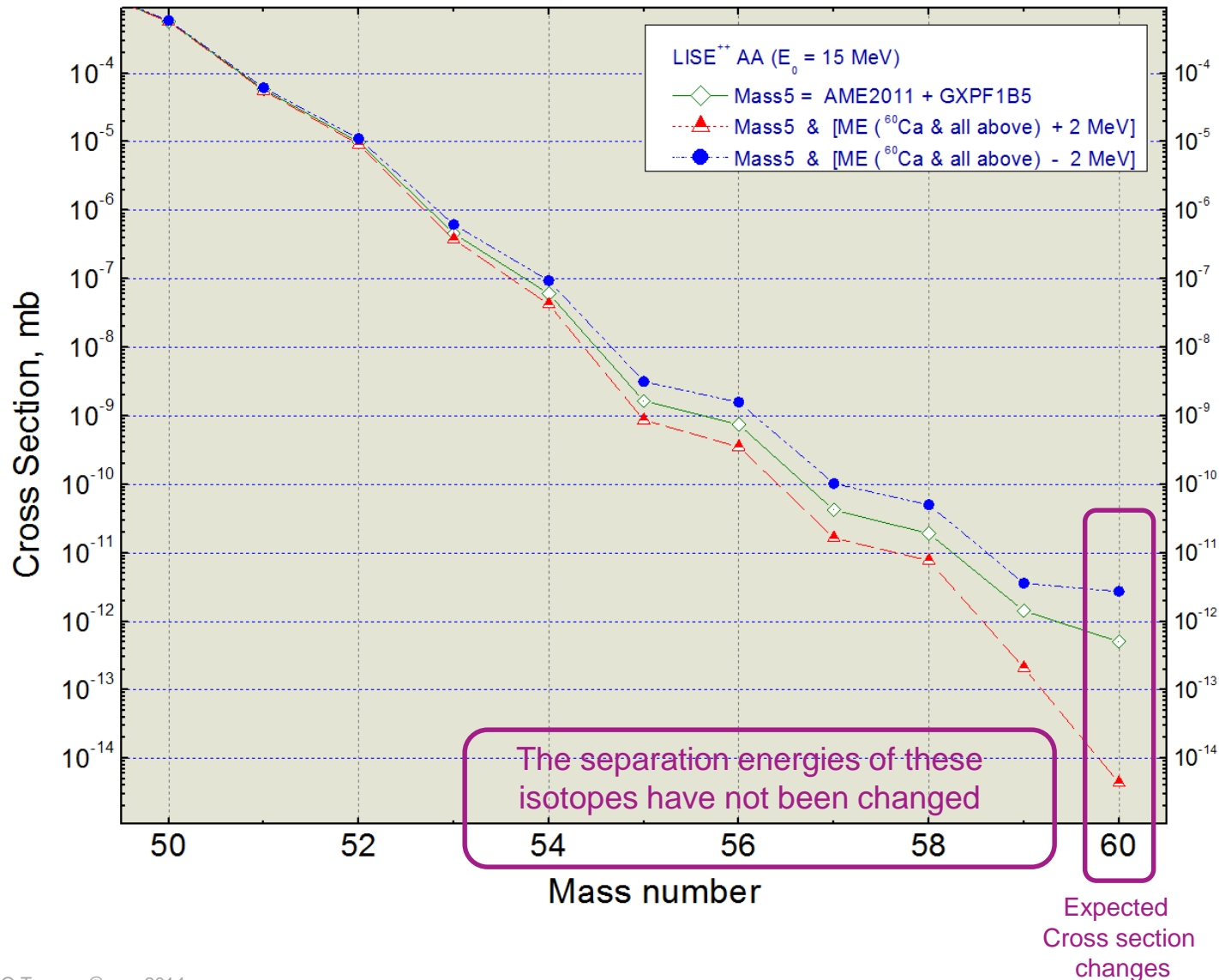
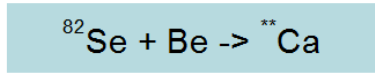
$^{238}\text{U}$  beam:  
 $\langle E \rangle = 27$  MeV (K.H.S.)

$^{82}\text{Se} + \text{Be} \rightarrow \text{Ca}^{**}$



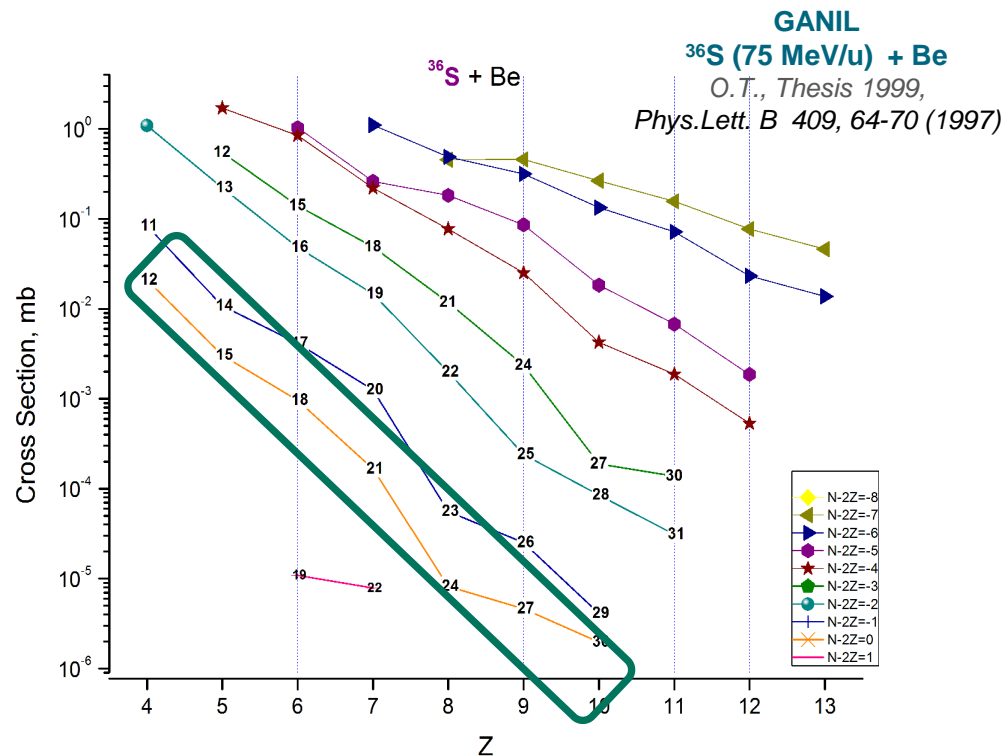
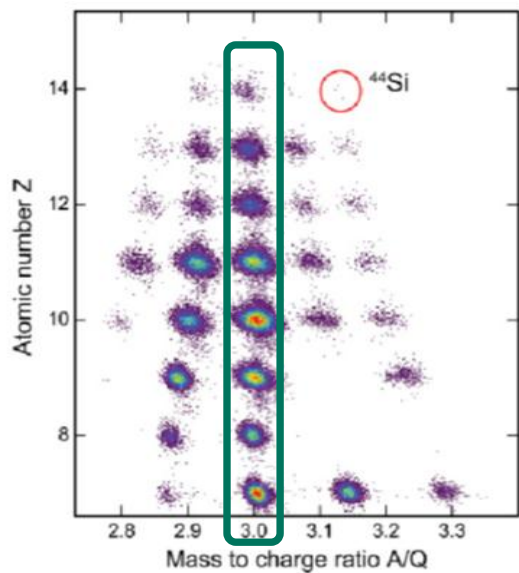
Let's **increase** (decrease)  $^{60}\text{Ca}$  and heavier isotopes mass excesses by 2 MeV  
 (or **decreased** [increased] neutron separation energies)





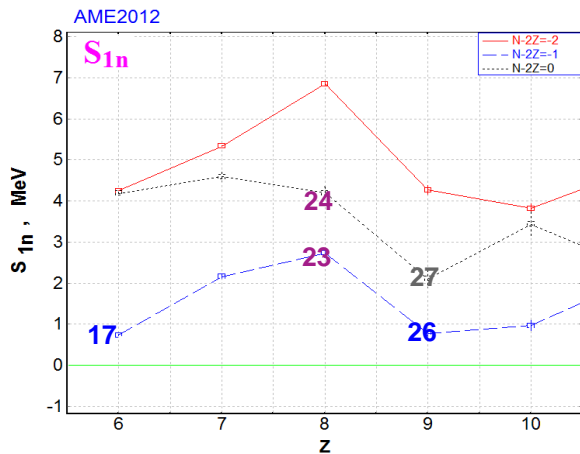
1. Separation energy changes influence drastically of cross section close to the drip-line (this  $^{60}\text{Ca}$  example)
2. Residue cross section depends how much bound are preceding isotopes
3. Deviations in cross sections are not just indicators for local low separation energies (as  $^{31}\text{Ne}$ ,  $^{37}\text{Mg}$ ), but also might provide information about shell effects close to the neutron drip-line

NSCL/MSU  
 $^{48}\text{Ca}$  (140 MeV/u) + Ta  
 O.T. et al., Phys.Rev. C 75, 064613 (2007)

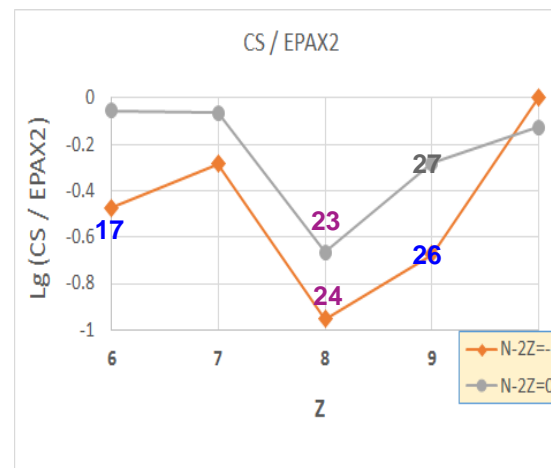


GANIL  
 $^{36}\text{S}$  (75 MeV/u) + Be  
 O.T., Thesis 1999,  
 Phys.Lett. B 409, 64-70 (1997)

Oxygen isotopes are more particle bound

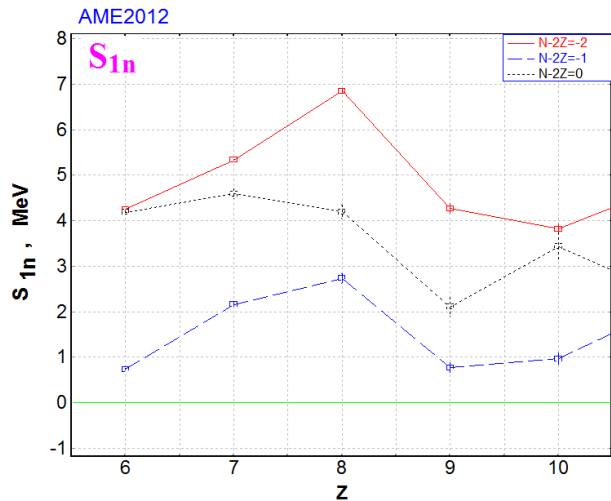


Oxygen isotopes are less produced

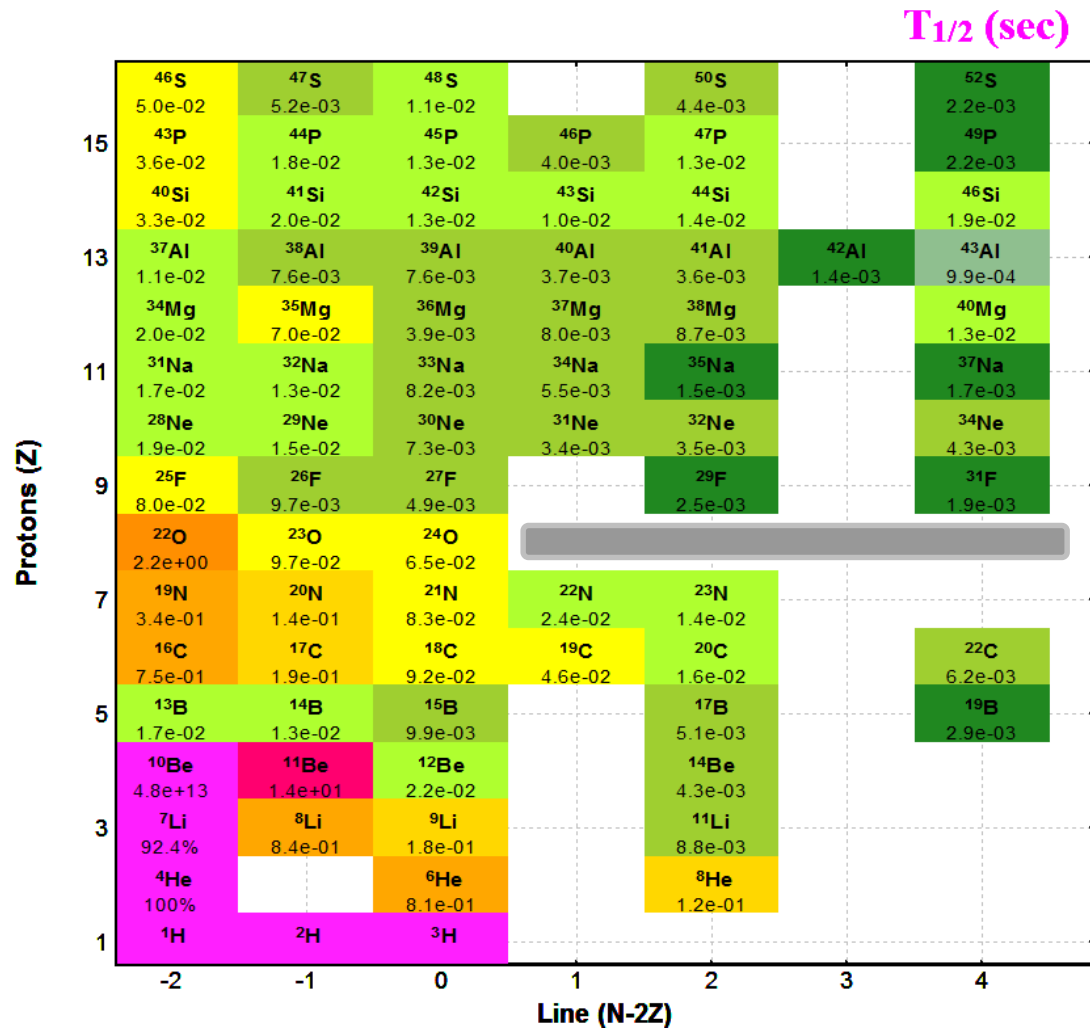
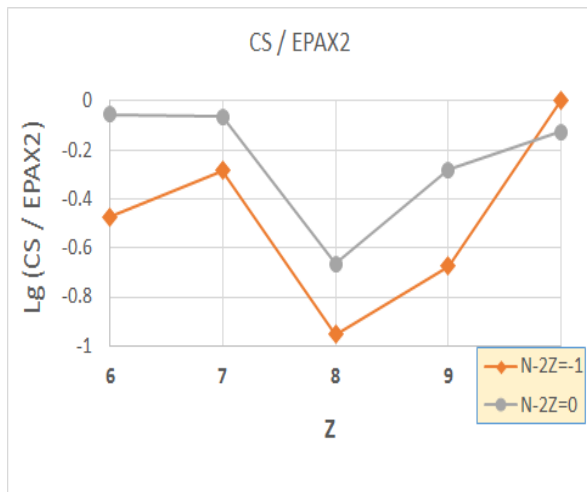


?

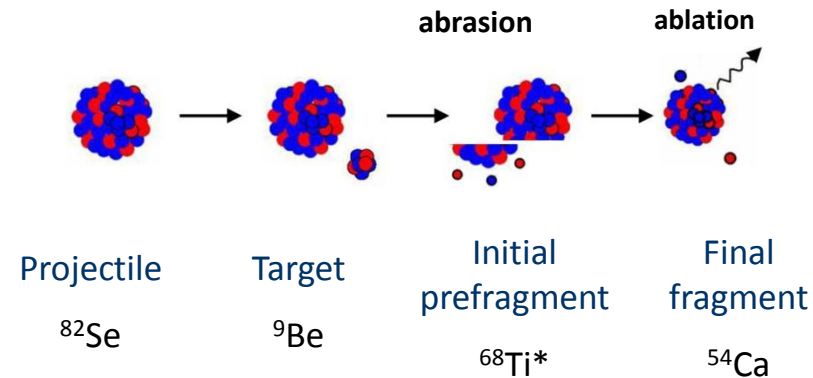




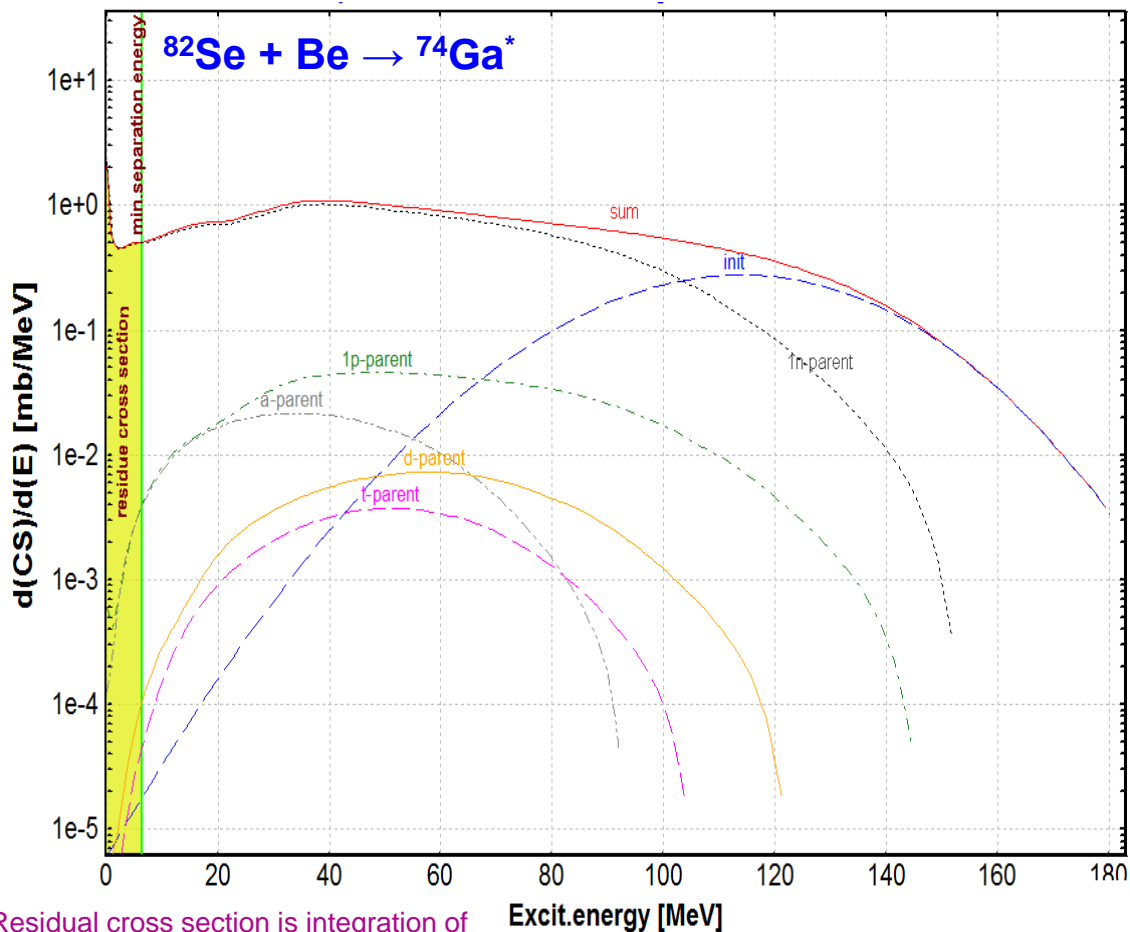
Oxygen isotopes are more particle bound, but less produced !?



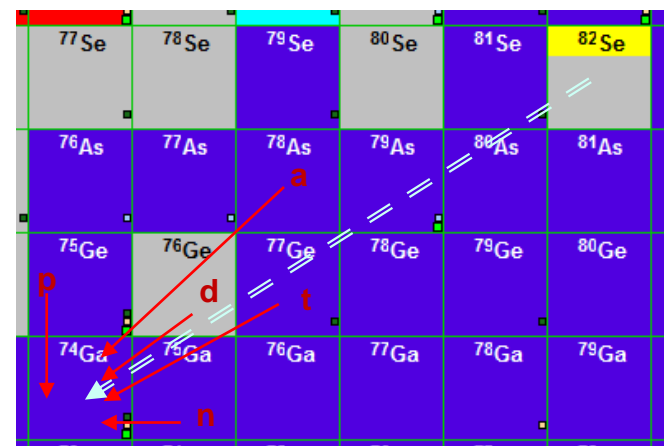
- No particle bound preceding isotopes of the same element, So "excitation energy train" cannot be slow down
- Absence of excited bound states?



- ❑ The LISE<sup>++</sup> AA model is initially based on the version of J.-J.Gaimard and K.H.Schmidt, NPA531 (1991) 709
- ❑ The LISE<sup>++</sup> AA model is analytical, that allows to calculate low cross sections of very exotic nuclei
- ❑ The Abrasion-Ablation approach meets three principal difficulties
  - a. Determination of Excitation energy parameters (models) for each reaction
  - b. Plenty of other parameters
  - c. Suggesting negligible contribution of dissipation processes during abrasion (it can be true at high energies with light targets)
- ❑ Four excitation energy models are implanted in the code
- ❑ The Ablation step (Evaporation cascade) uses a mass table to obtain separation energies



Residual cross section is integration of the excitation function from 0 to the minimum of separation energy

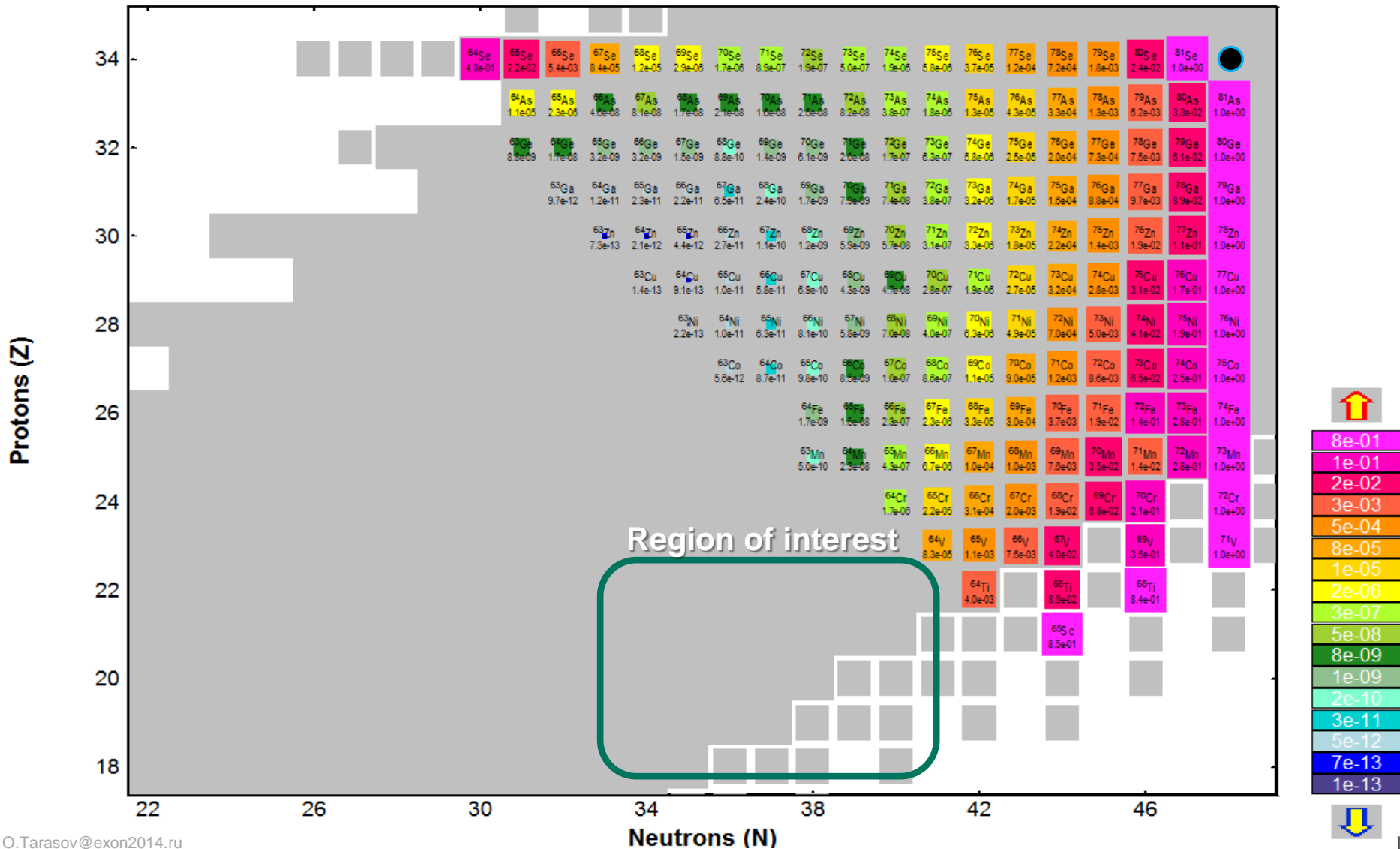


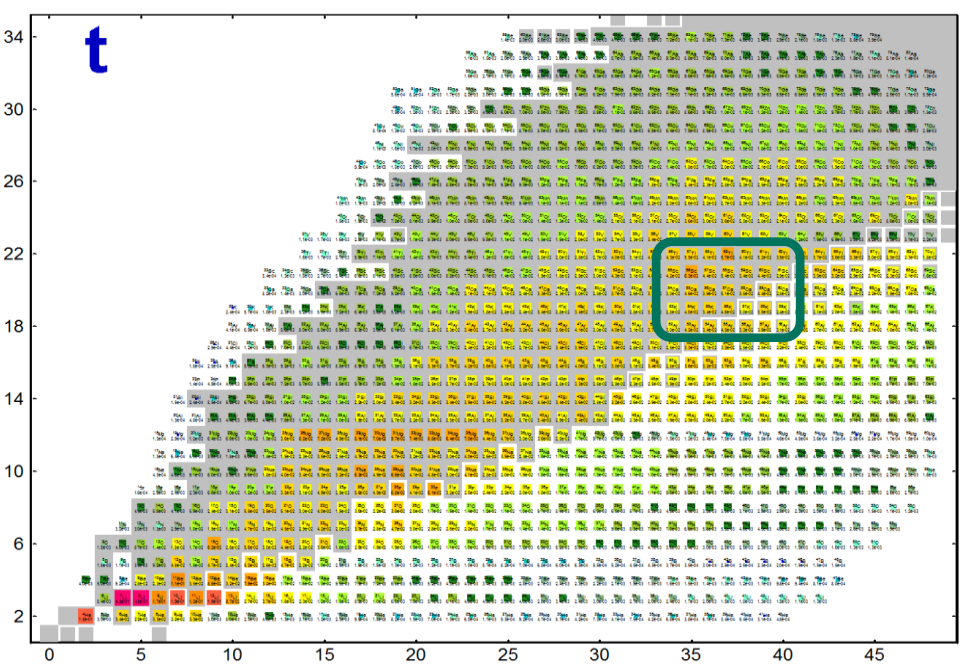
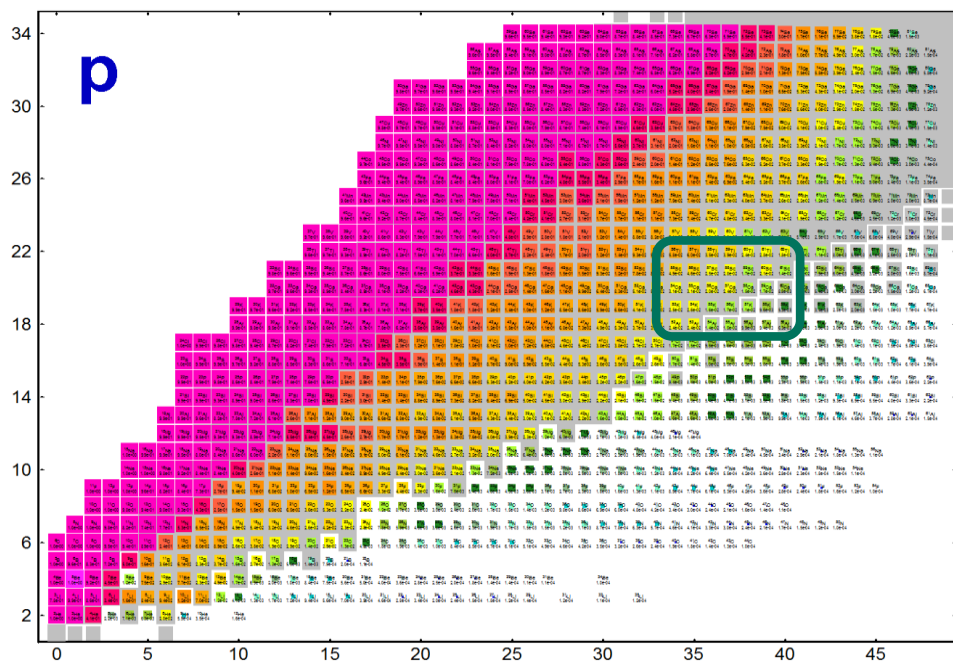
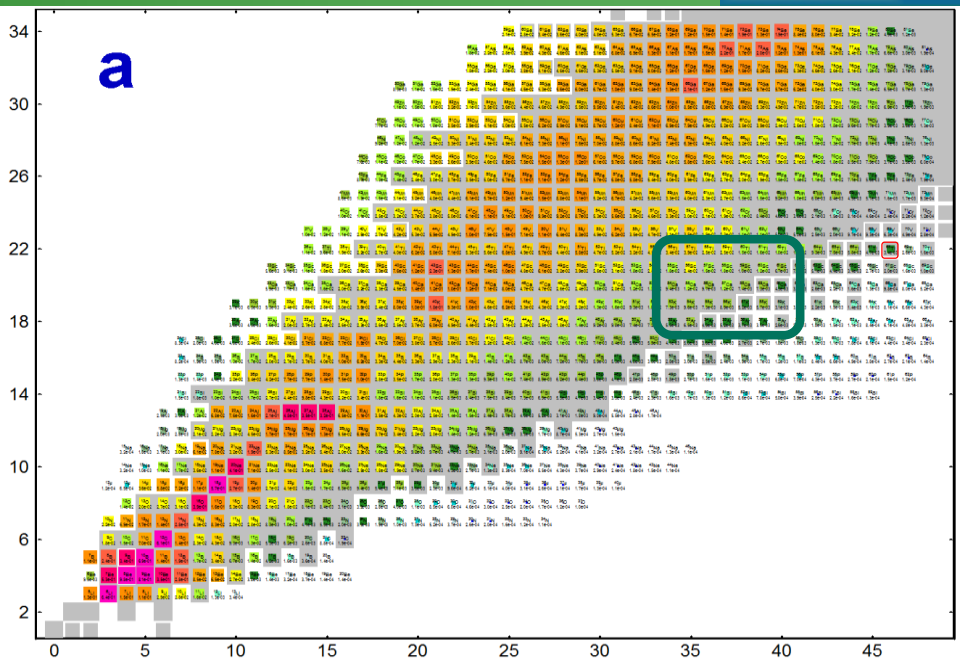
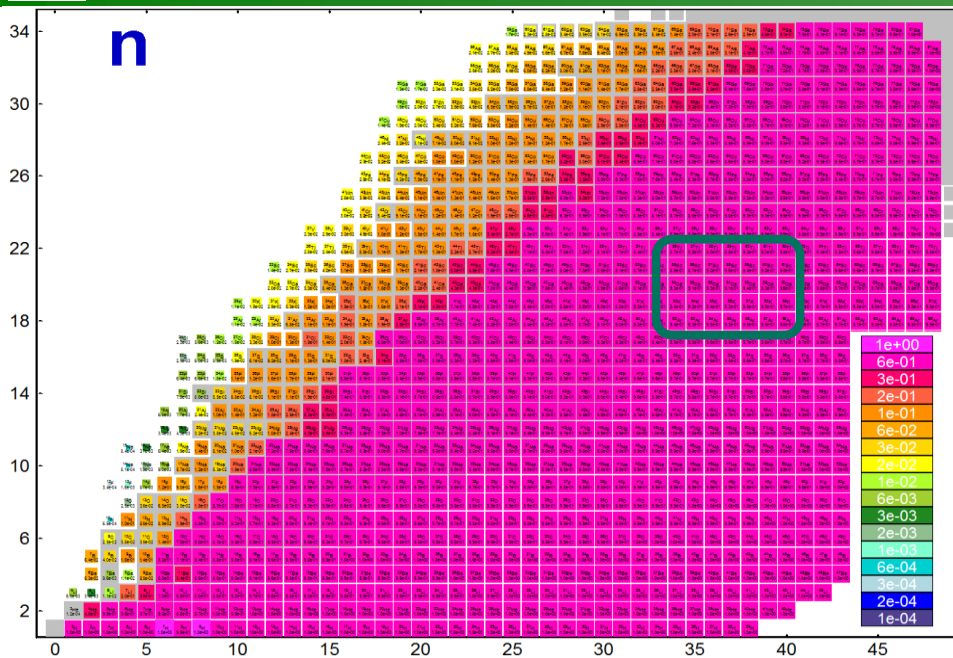
1. Largest incoming contribution to the Total excitation function is 1n-channel
2. Largest incoming contribution to the Residue cross section is 1n-channel

**Current mode: Initial CS -> [ S residue ] / [ Sr total ]**

ABRASION-ABLATION -  $^{82}\text{Se} + \text{Be}$

Excit.Energy Method:< 2 >; < E\* >:15.0\*dA MeV Sigma:9.15; Coef<sup>Thermalization</sup>=5.00e-22 MeV.s DB<sub>1</sub>="GXPF1B"  
 NP=64; SE:"DB1+Cal0" Density:"auto" GeomCor:"On" Tunlg:"auto" FisBar=#1 Bar<sup>Fac</sup>=1.00 Modes=1010 1010 010

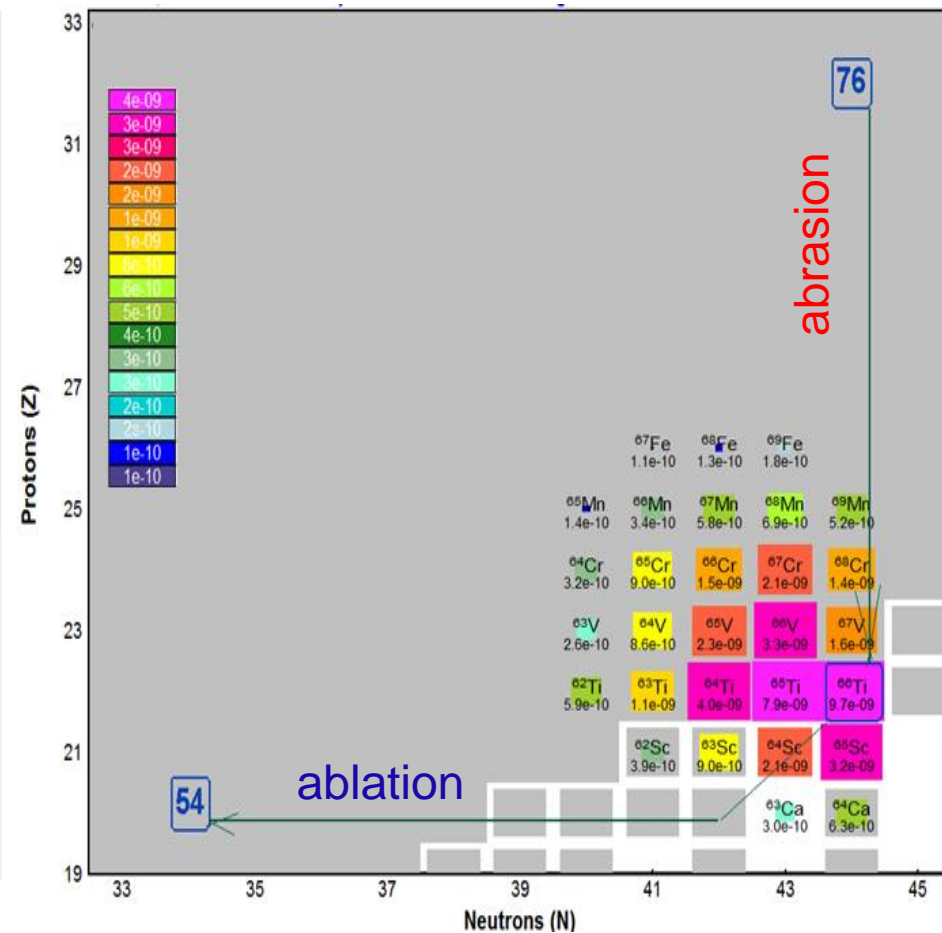
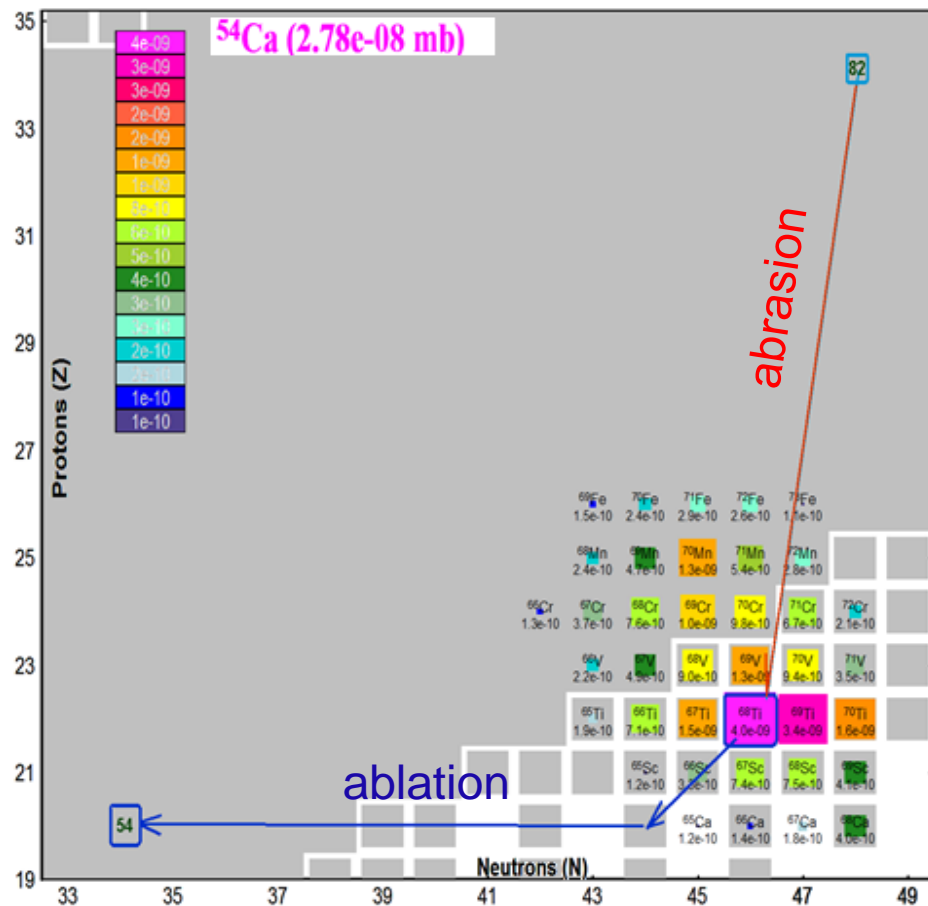




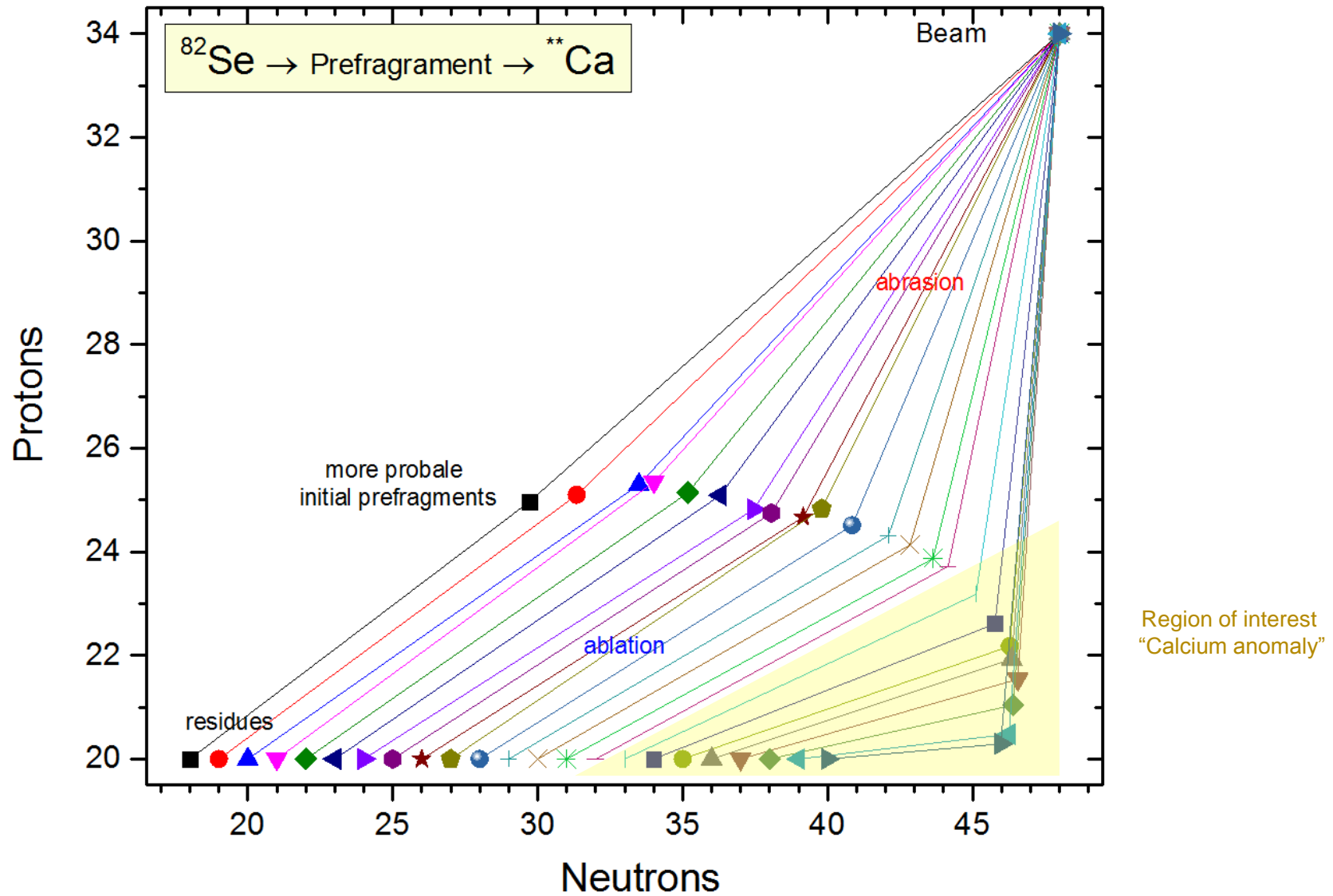


<sup>82</sup>Se + Be

<sup>76</sup>Ge + Be

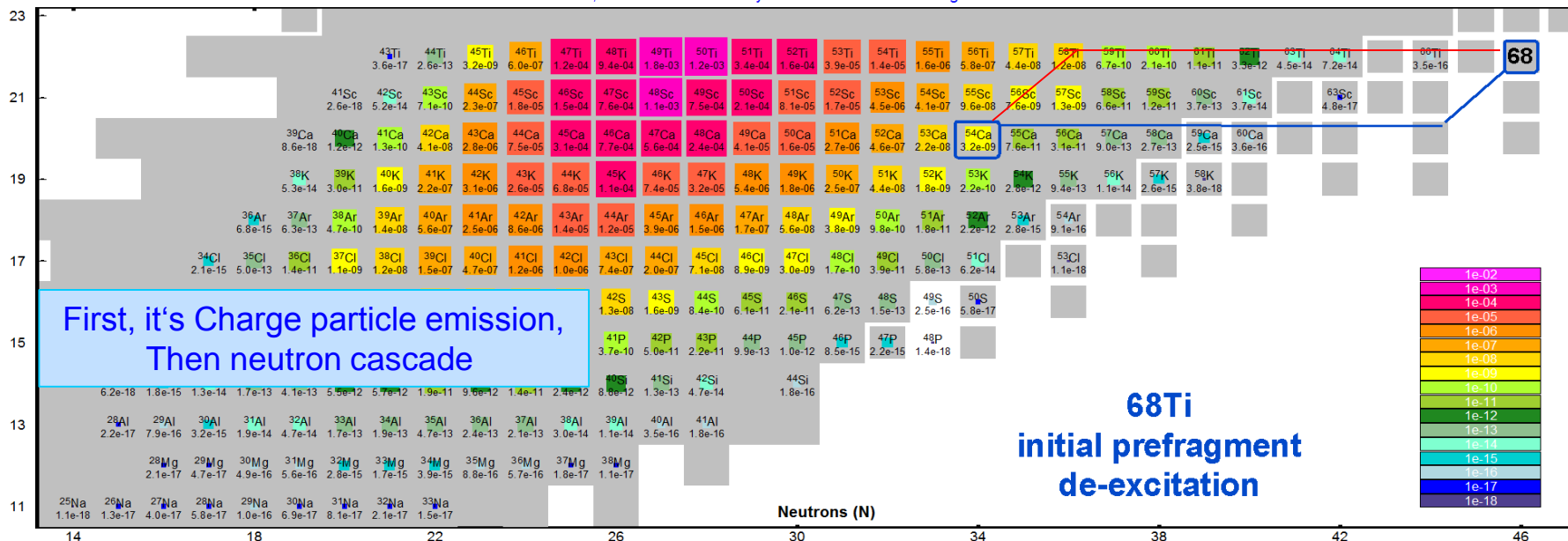


More probable prefragments are Ti-isotopes (dZ=2)



## Final Evaporation Residue cross-sections (LisFus)

EVAPORATION - Compound nucleus  $^{68}\text{Ti}$   
 Excit.Energy: 149.0-207.0 MeV; Fus.CS: 0.0 mb; Fus.Barrier: 10.82 fm;  $h_\omega = 2.0$  MeV  
 NP=64; SE:"DB1+Cal0" Density:"auto" GeomCor:"On" Tunlg:"auto" FisBar=#1 Bar<sup>Fac</sup>=1.00 Modes=1010 1010 010



First, it's Charge particle emission,  
Then neutron cascade

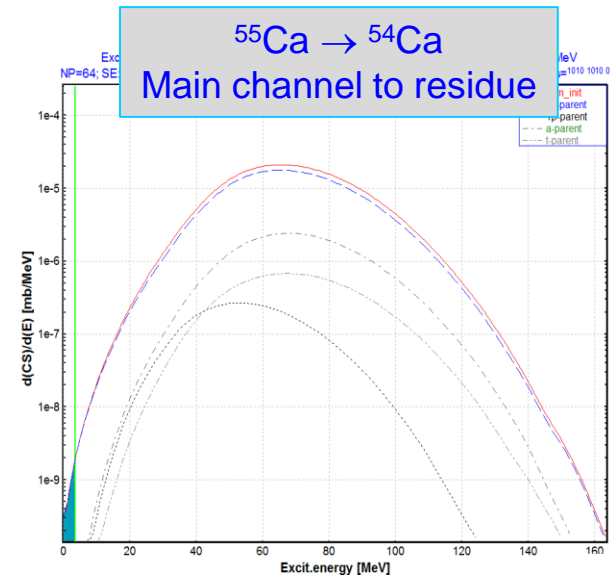
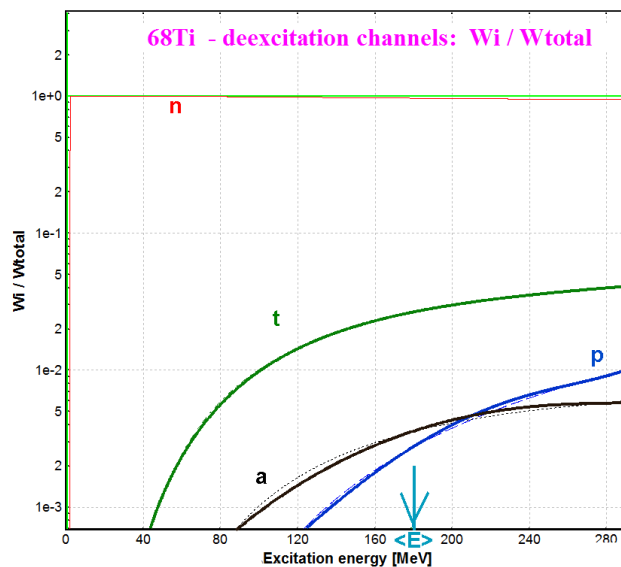
Probability for  $^{68}\text{Ti}^*(\text{Ex}=180\text{MeV})$

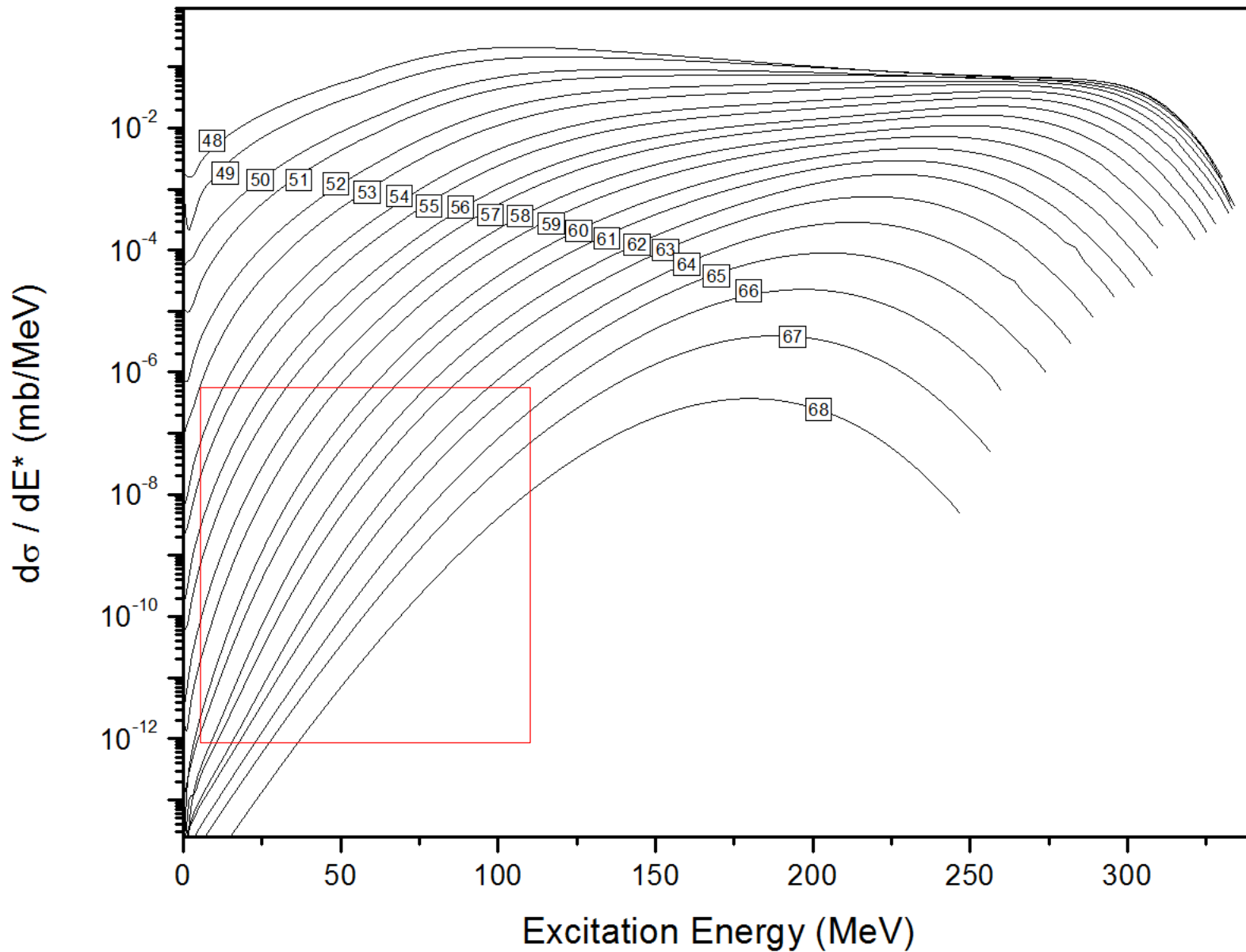
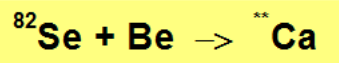
- $t = 2.6\text{e-}2$
- $a = 3.6\text{e-}3$
- $p = 9.3\text{e-}3$

Probability ( $dZ=2$ )

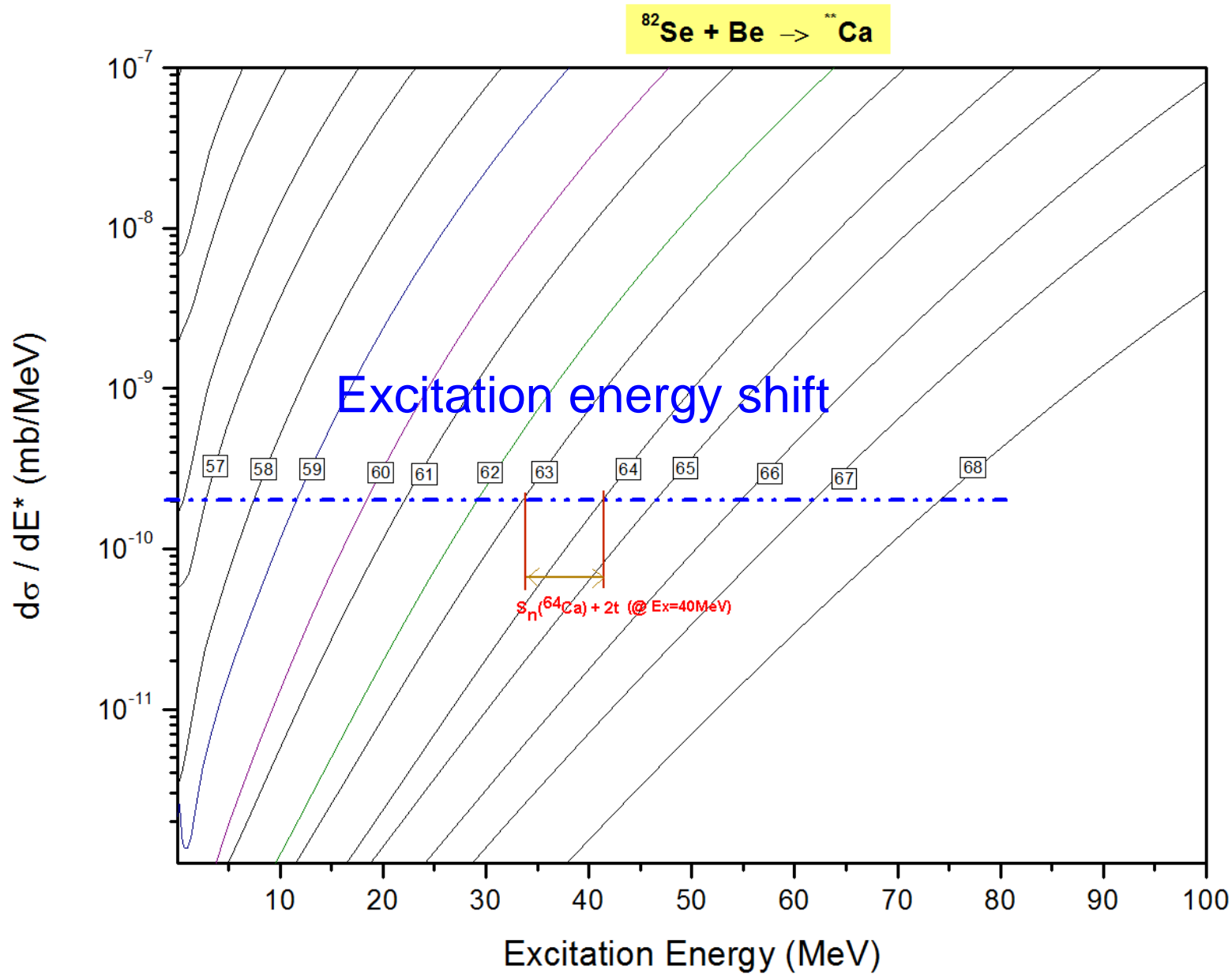
- $t^2 = 6.8\text{e-}4$
- $a = 3.6\text{e-}3$
- $p^2 = 8.7\text{e-}5$

It is necessary to create the MC version to gate for  $^{54}\text{Ca}$  residual in order to answer where ( $^{68}\text{Ti} \rightarrow ^{54}\text{Ca}$ ) de-excitation by charge particles is more probable

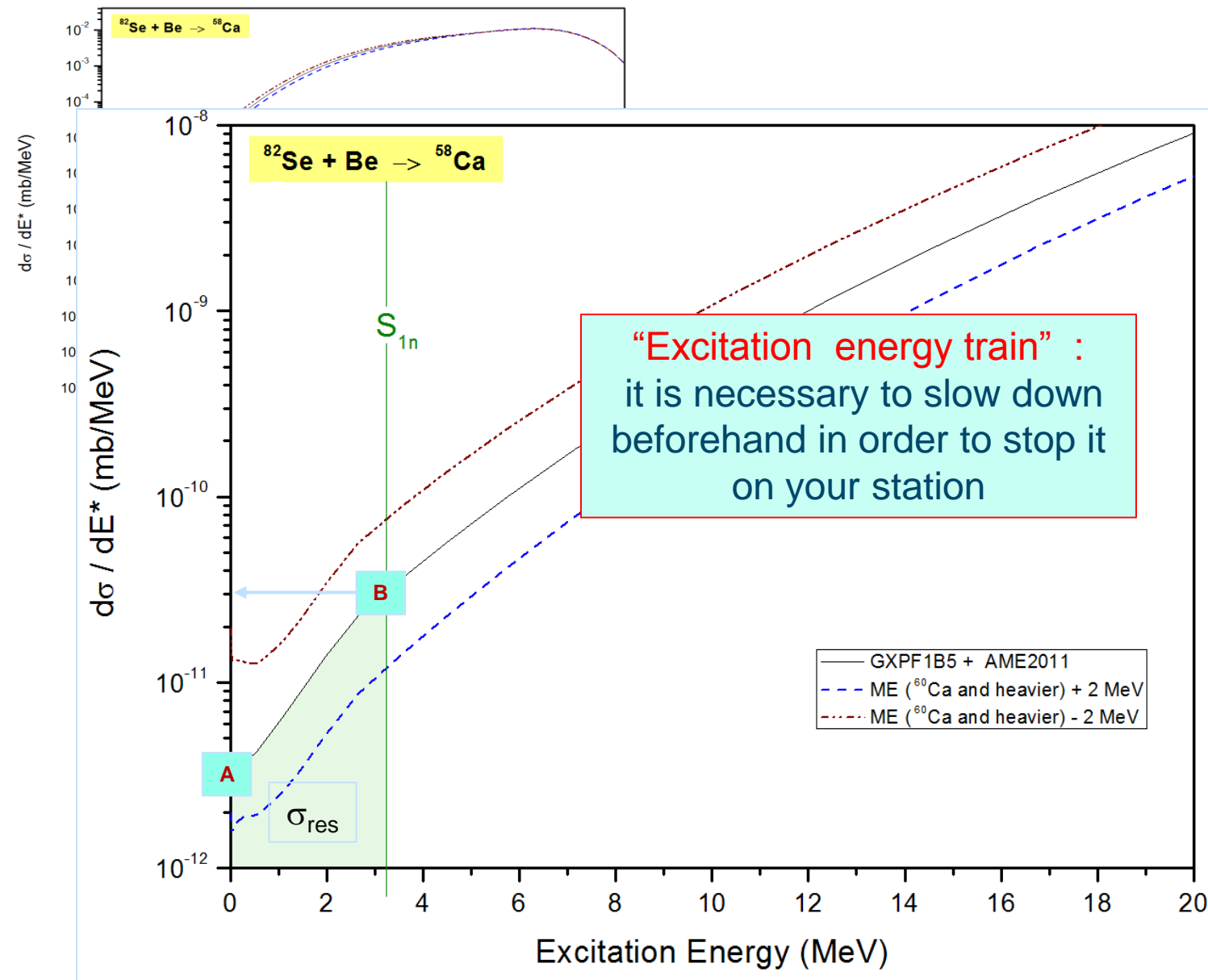




Zoom







$${}^{58}\text{Ca } \sigma(0) = A$$

$${}^{58}\text{Ca } \sigma(3.2) = B$$

1.  $S_{1n}({}^{58}\text{Ca}) = 3.2 \text{ MeV}$  :  
 $\sigma(S_{1n}) = B$  is moving to  
 $E^* = 0$  for  ${}^{57}\text{Ca}$  excitation  
 function  
 ${}^{57}\text{Ca } \sigma(0) = B$

2. Let's assume  
 $S_{1n}({}^{58}\text{Ca}) = 0 \text{ MeV}$ , then  
 in the first rough  
 approach :

$\sigma(0) = A$  is the same  
 value @  $E^* = 0$  for  ${}^{57}\text{Ca}$

$${}^{57}\text{Ca } \sigma(0) \sim A$$

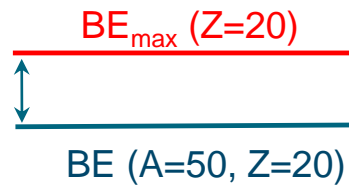
$${}^{57}\text{Ca } \sigma(2.6) \sim B$$

$$\Sigma S_n(A,Z) = BE_{\max}(Z) - BE(A,Z)$$

Lets call dBE as

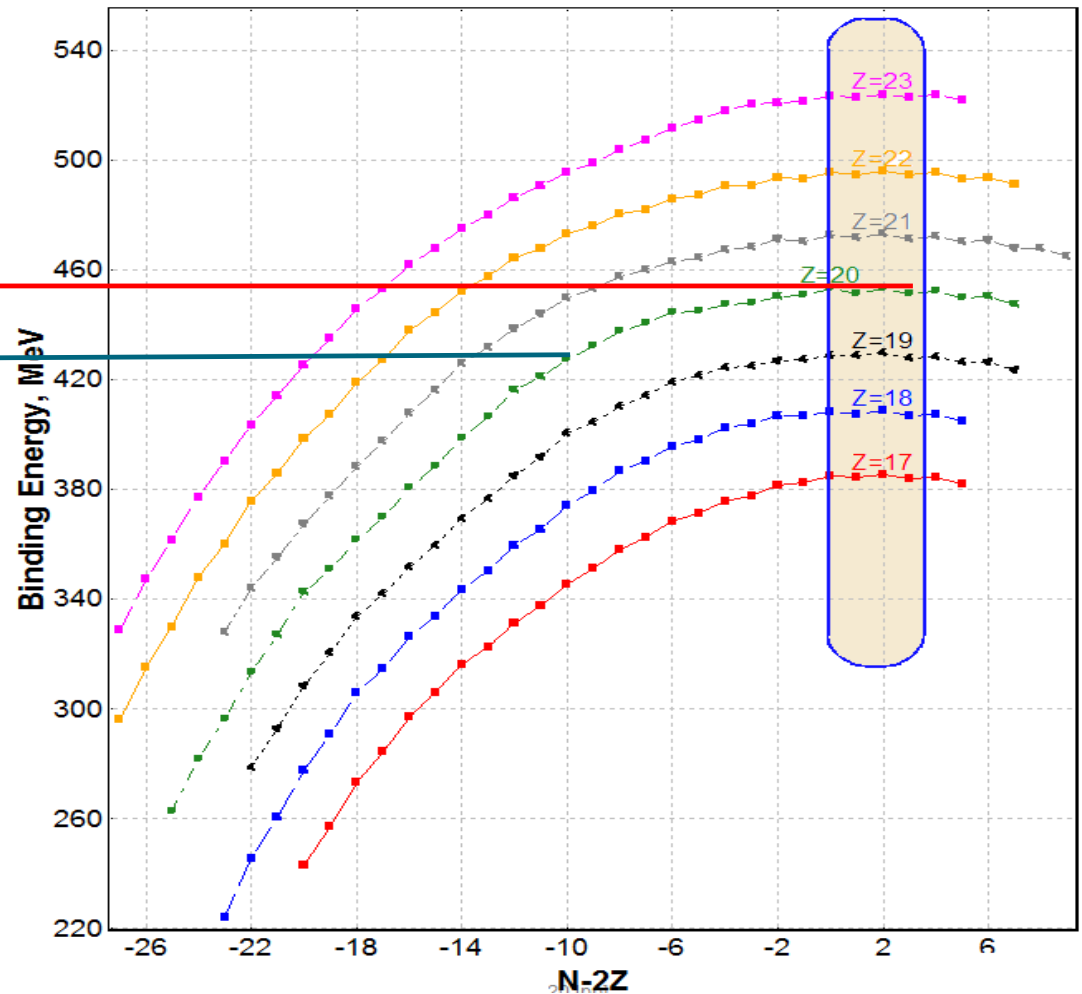
$$dBE(A,Z) = BE_{\max}(Z) - BE(A,Z)$$

dBE (A=50,Z=20)

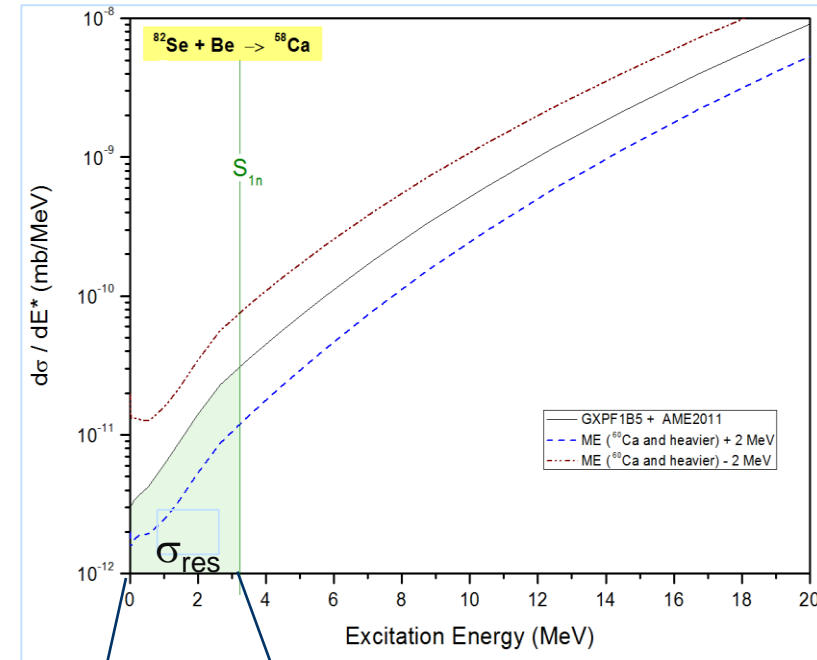
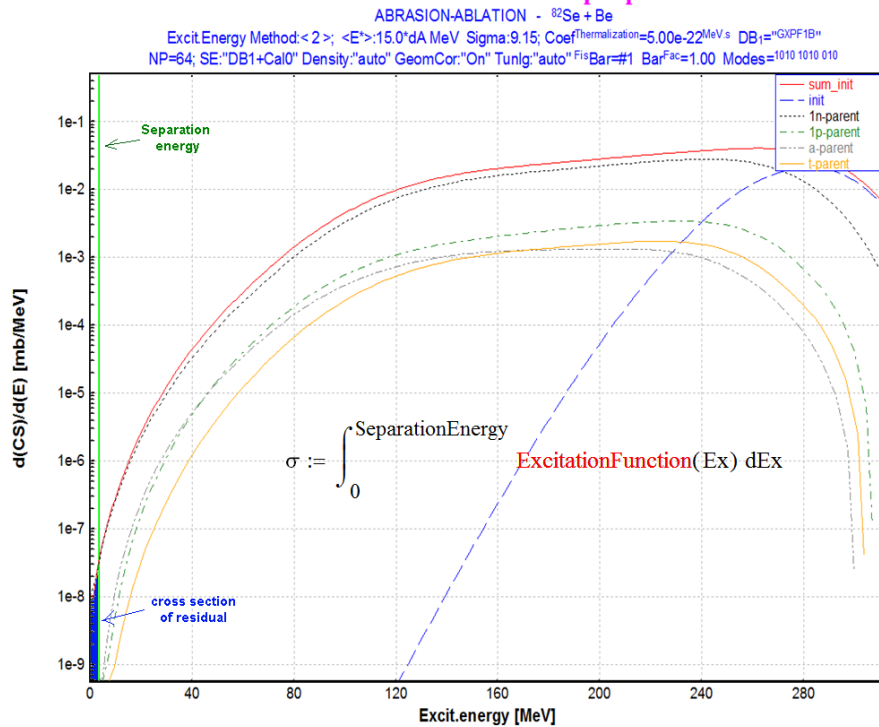


## Binding Energy

<Database: User's ME file (GXPF1B) + LDM0>  
Z=17-23



## 54Ca excitation distributions: Input parent distr.



$dBE(Z,N)$

$dBE(Z,N+1)$

$$\sigma(Z,N) = f [ dBE(Z,N), S_n(Z,N) ]$$

$S_n(Z,N)$  minimum separation energy

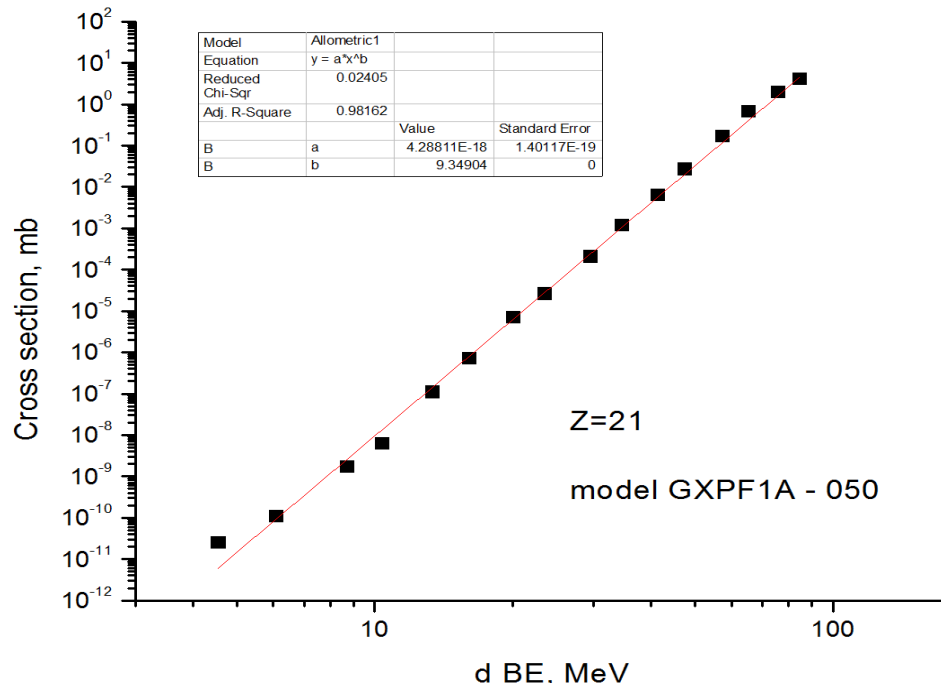
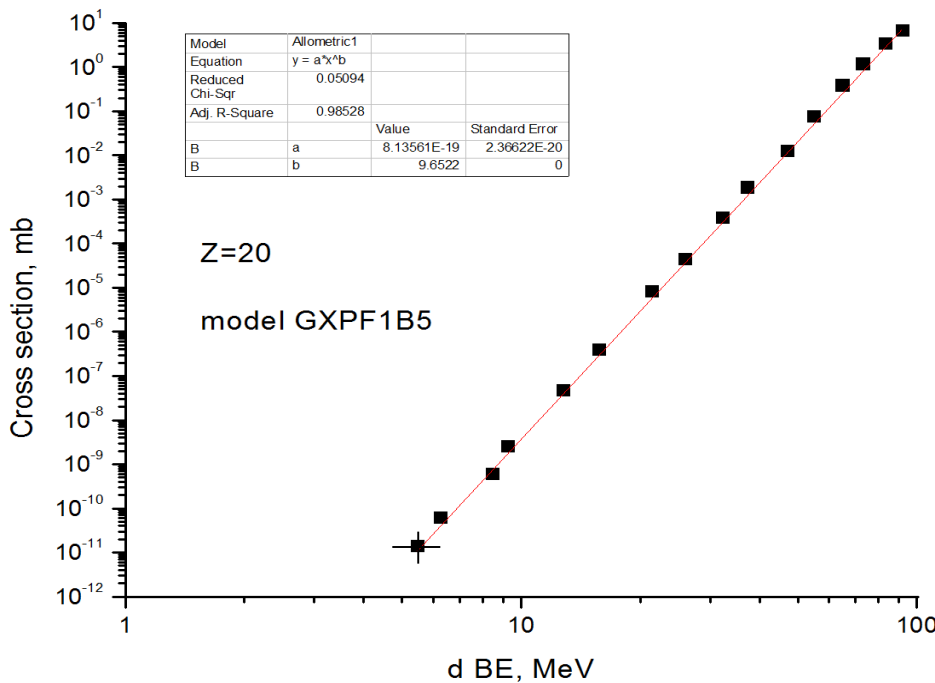
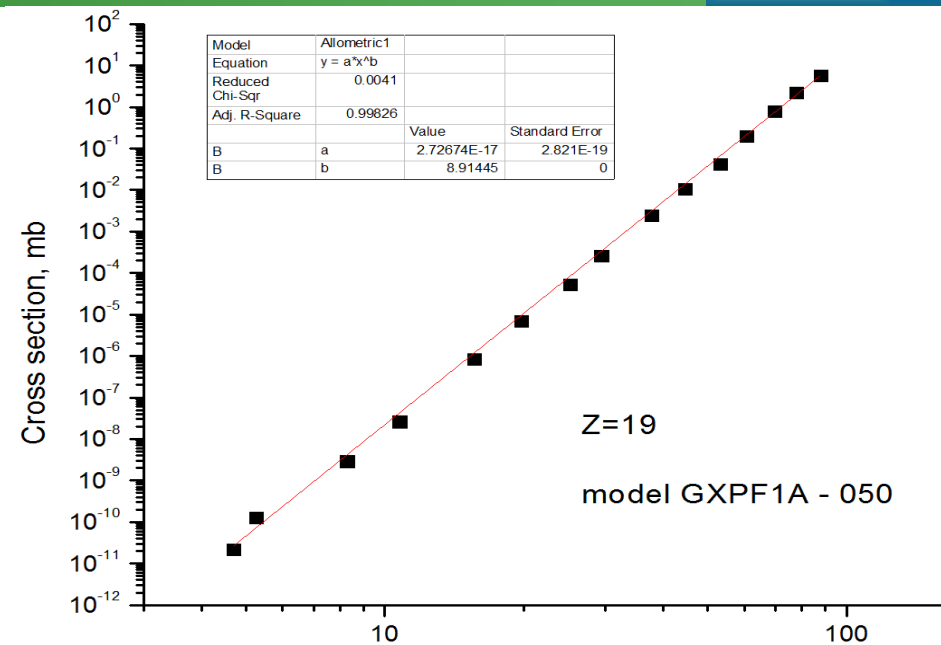
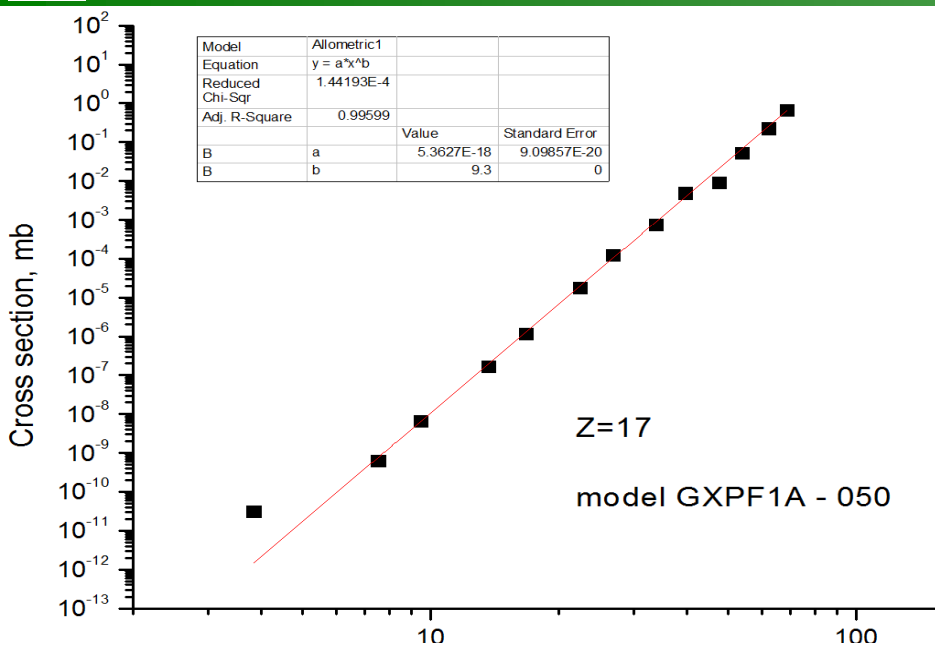
$dBE(Z,N)$  difference between the maximum binding energy for isotopes (Z) and binding energy of the nucleus (Z,N)

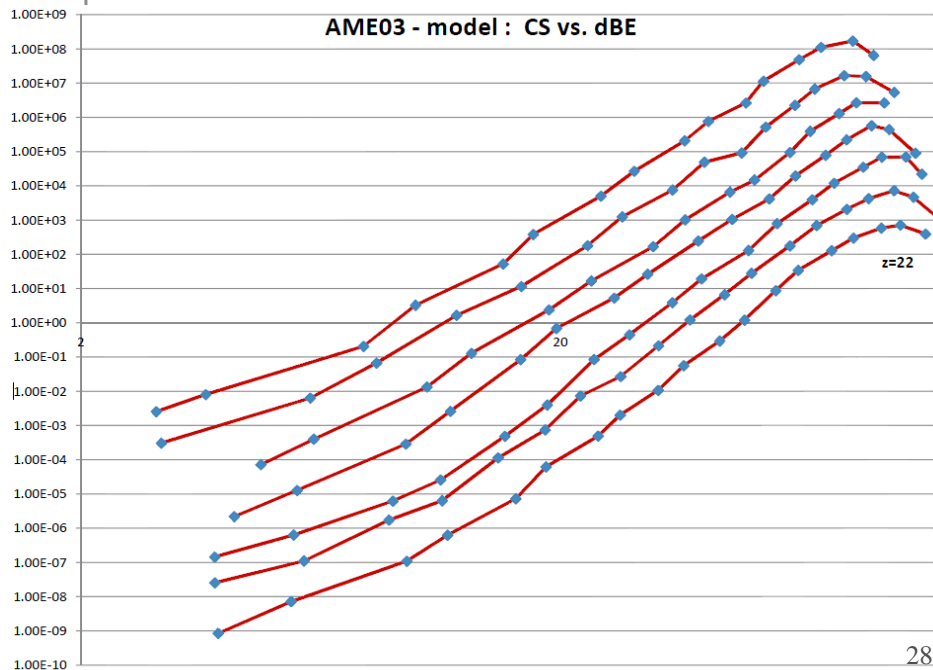
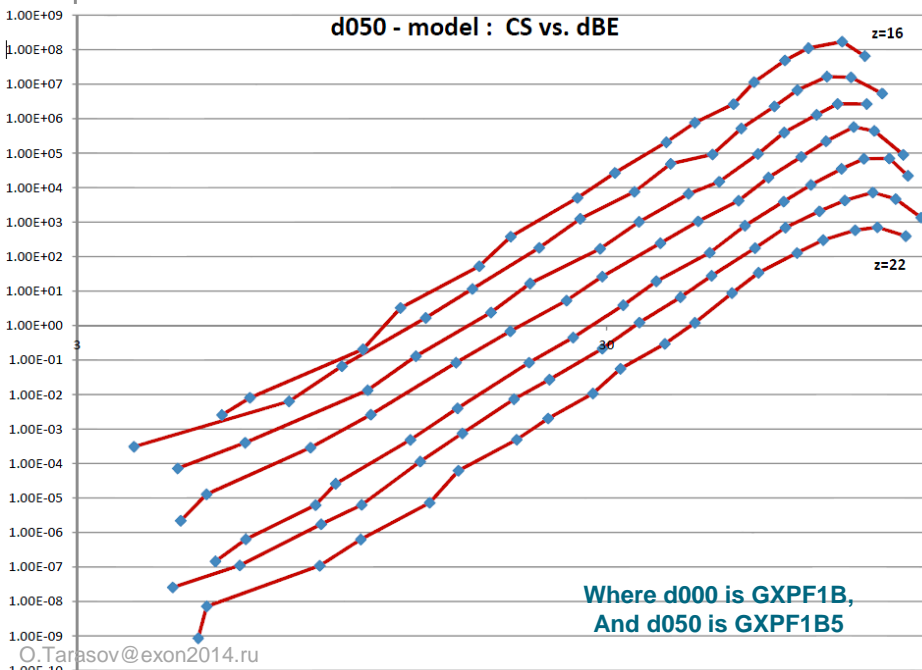
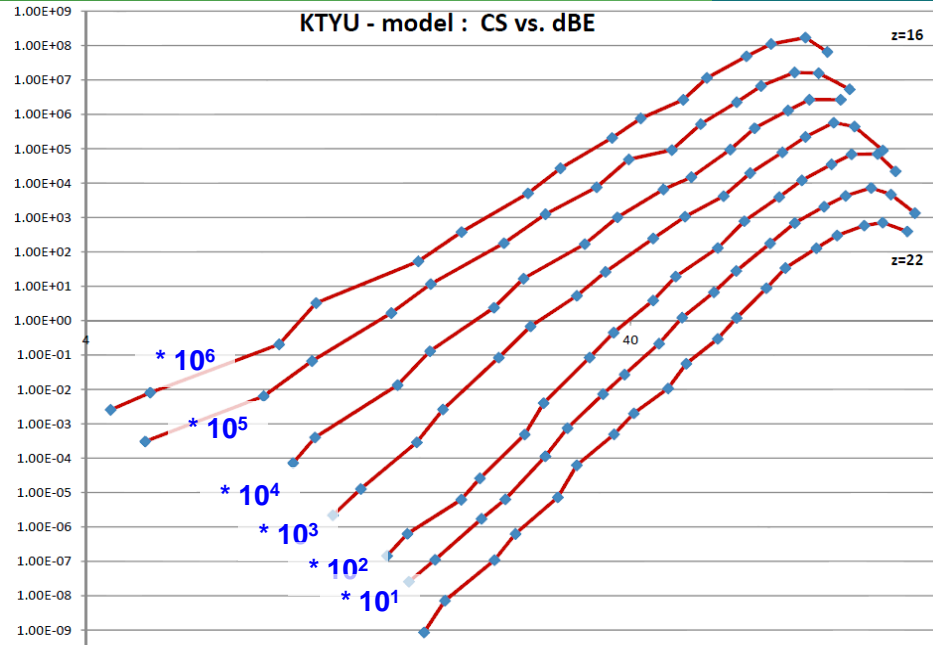
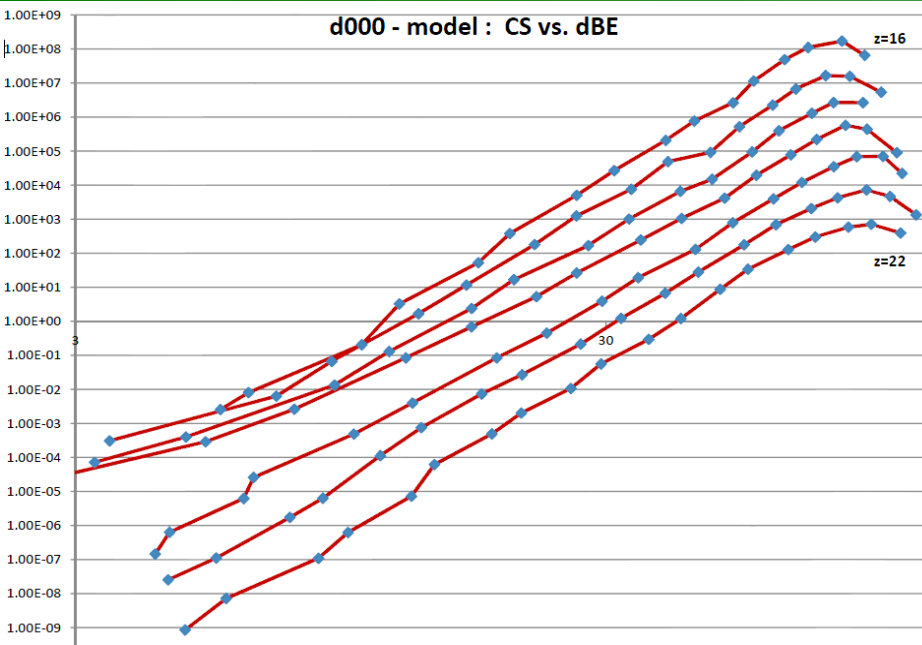
$$1. \quad \sigma(Z,N) \sim [dBE(Z,N) + a_1]^{a_2} * S_n(Z,N)$$

If  $S_n(Z,N) \leq 0$ , Then  $\sigma(Z,N) = 0$ ,  
 whereas  $Q_\alpha$  or  $BE/A$  systematics show unbound nuclei

Or Using  $dBE(Z,N-1) = dBE(Z,N) + S_n(Z,N)$ ,

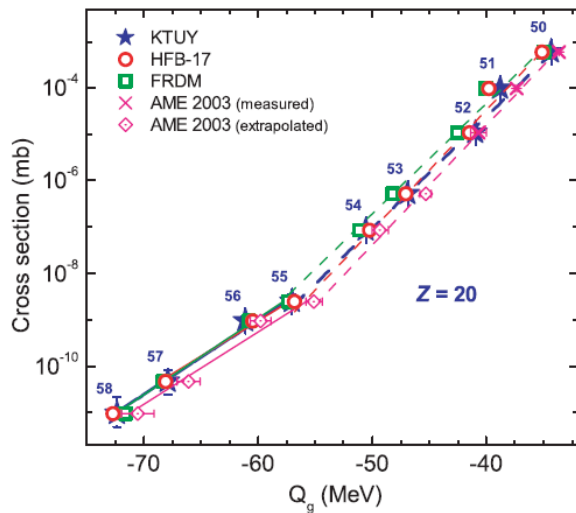
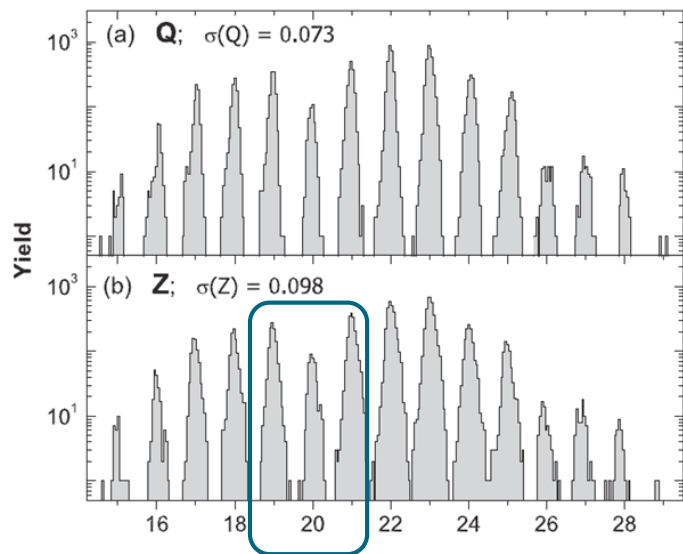
$$2. \quad \sigma(Z,N) \sim dBE(Z,N+1)$$



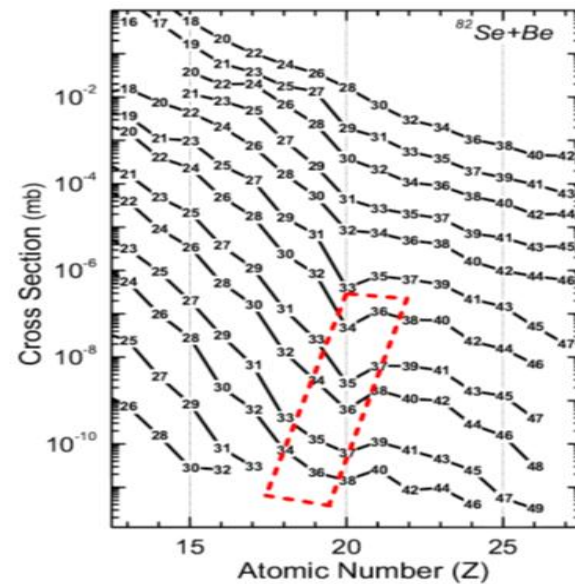
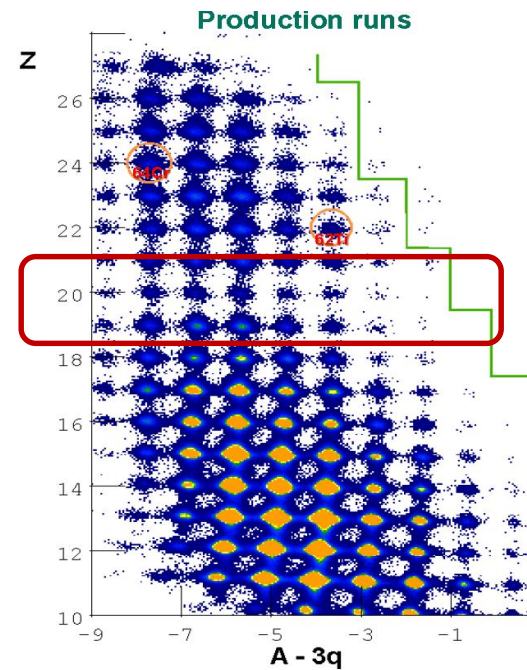


Where d000 is GXPF1B,  
And d050 is GXPF1B5

<sup>76</sup>Ge



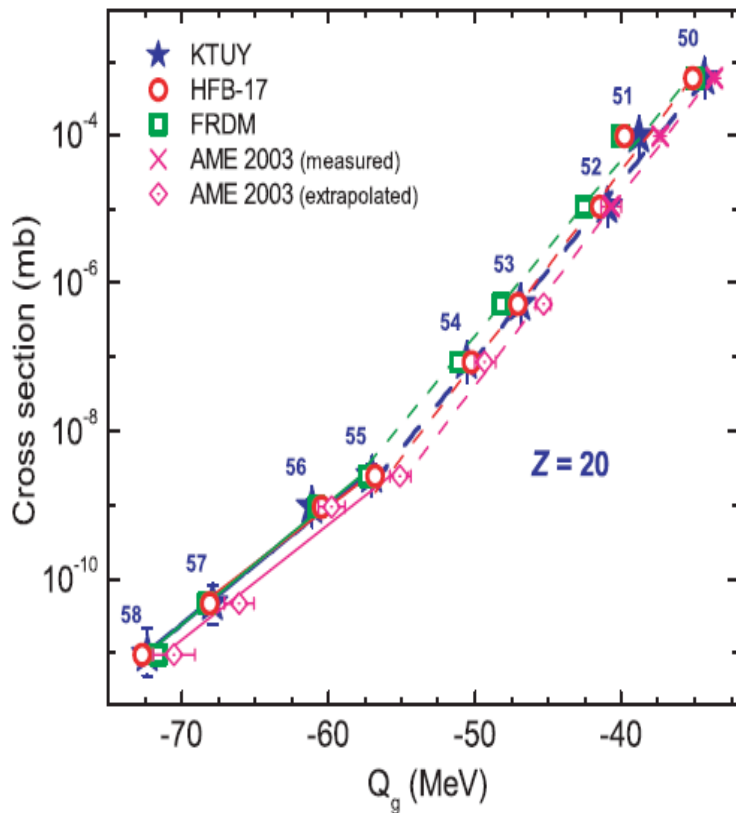
<sup>82</sup>Se







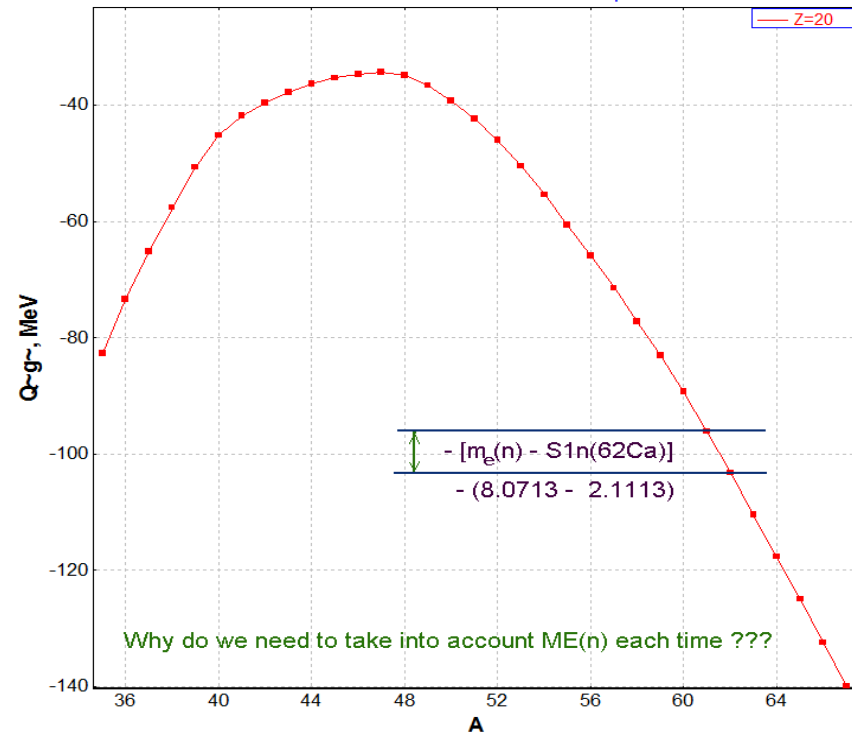
$^{82}\text{Se}$



## $Q_g$ systematics (P $\rightarrow$ F)

$^{82}\text{Se} \rightarrow Z=20$

Odd-even corrections : neutrons - Yes; protons - No



In  $Q_{gg}$  systematics “this” neutron is compensated by conjugated products

1. Even one registered event (non-zero production cross section) provides information for nucleus structure : particle bound or not
2. There are very low statistics, and spectroscopy experiments are difficult (or impossible) to perform, Systematic production cross sections can provide some indications about structure of observed isotopes, and even provide hints about structures of preceding non-observed isotopes
3. Even it is difficult to evaluate masses from AA analysis (one particular cross section kink, or depression for several isotopes), though production cross section analysis is powerful test of theoretical mass models
4. dBE-systematics advantages:
  - a. deduced, not assumed
  - b. start energy point
  - c. unbound nuclei are out of the systematics
  - d. can be used for other reaction mechanisms, where neutron rich nuclei are produced after emission large number of neutrons
  - e. no parameters, the same slope?
  - f. no any odd-even corrections and so on.

1. What is proper expression (parameterization) for dBE-systematics?
2. Why slopes are similar for all elements?
3. Why is turnover point is always in the dBE-place (stability line?)?
4. What is contribution of dissipation at low energies or with heavy targets
5. Can Monte Carlo study of prefragment distribution and channels preceding to final residue explain the data?  
New experiments will be useful as well.

*Discussions with Prof. D.J.Morrissey  
are very appreciated.*