

Status after today's analysis: LISE⁺⁺ v18.5.14

Goal: understand and fix the low / too-narrow proton momentum distribution used by the Prompt Radiation utility.

Contents

- Momentum-model limitations for light charged particles: projectile-fragment models vs. de-excitation products.
- CM $\rightarrow$ Lab tests for $^{79}\text{Ge}^*$ and $^{60}\text{Ca}^*$: fixed CM angles, all-angle spectra, and TKE / MeV/u checks.
- PHITS-inspired reference plots are qualitative only; component contributions are not normalized.

Main result

- Problem found: 1H used $A_{pf} = 1$, giving $\text{sigparB} \approx 194$ and a direct-like, too-narrow distribution.
- Neighboring light particles used evaporation-parent-like $APF \approx 27\text{--}29$, giving $\text{sigparB} \approx 834\text{--}848$.
- Fix: SearchPrefragment() now uses a temporary 2H seed for requested 1H, while keeping final product = 1H.

Discussion / next checks

- Validate v18.5.14 spectra against requested PHITS/JQMD output from Tom before using absolute yields.
- Keep channel bookkeeping explicit: target-like evaporation, projectile-residue de-excitation, participant / pre-equilibrium.
- Recommendation: present v18.5.14 as a physics-motivated width correction, not only as a code patch.

Recommendation on the Convolution Model for Light-Particle Momentum Distributions in LISE⁺⁺

Main recommendation:

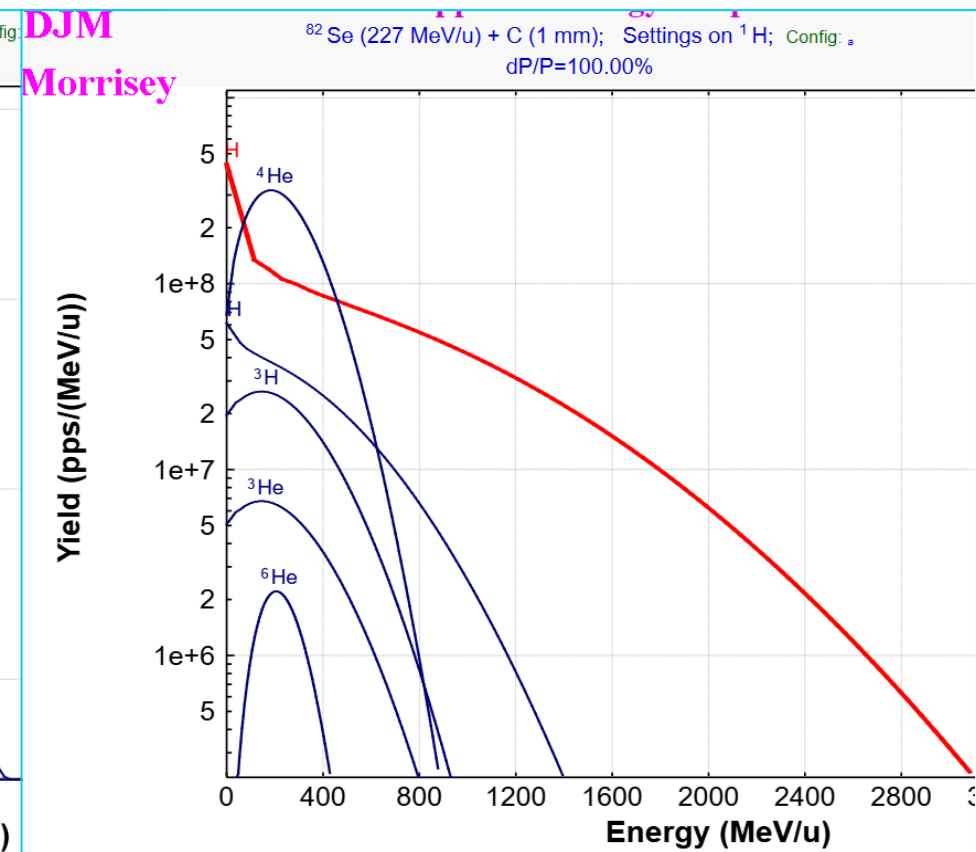
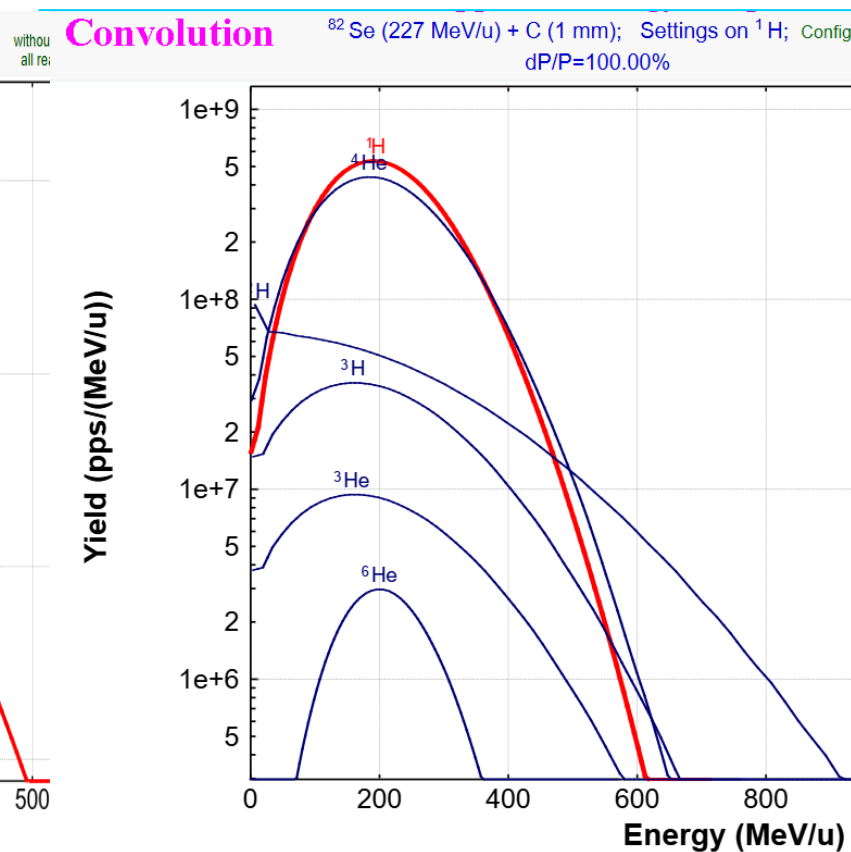
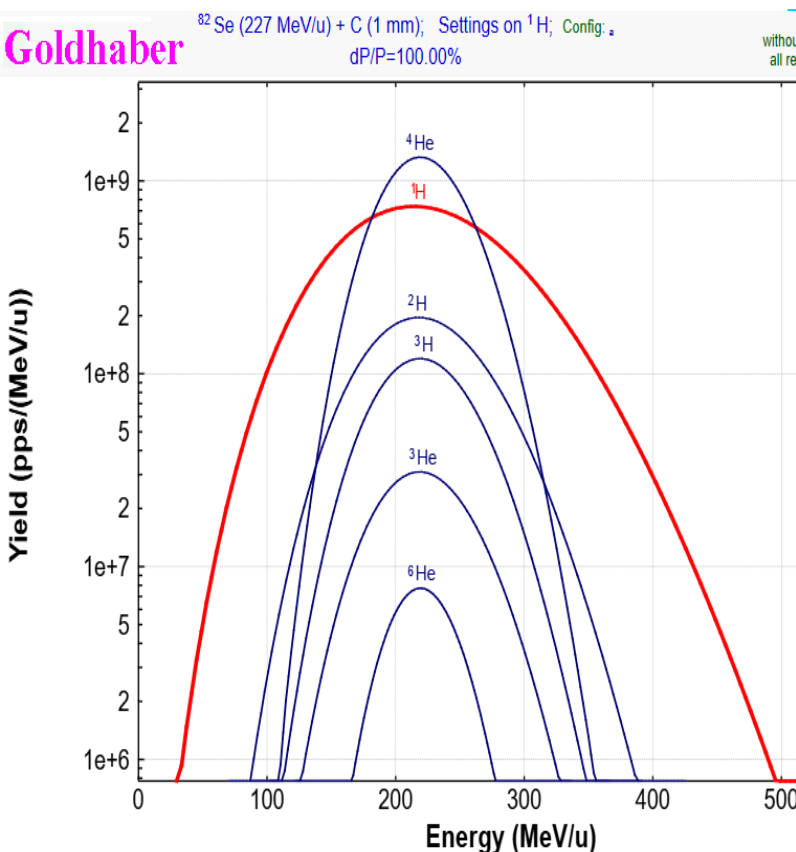
Develop and implement a new mechanism-dependent model for light-particle momentum and energy distributions in LISE⁺⁺. The existing models were developed mainly for projectile-fragmentation residues and do not explicitly include projectile de-excitation, participant/pre-equilibrium emission, target evaporation, breakup, and coalescence.

The Convolution model and the v18.5.14 correction improve the present calculations, but they represent only an interim solution.

Too narrow for de-excitation particles

1P = 4He .. Not so good

Too high production of 1H for 3 GeV



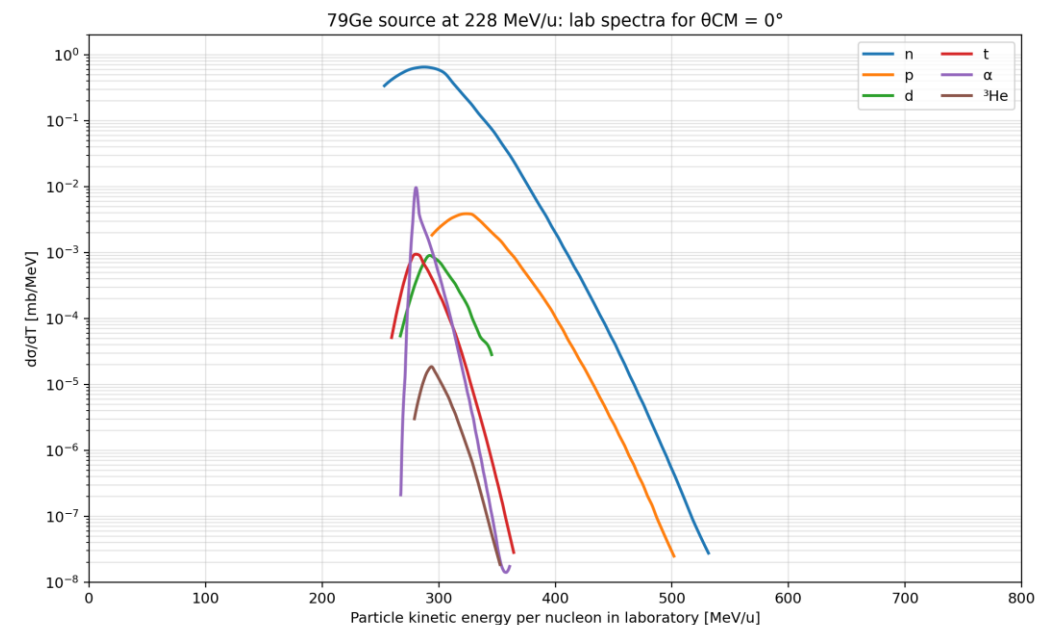
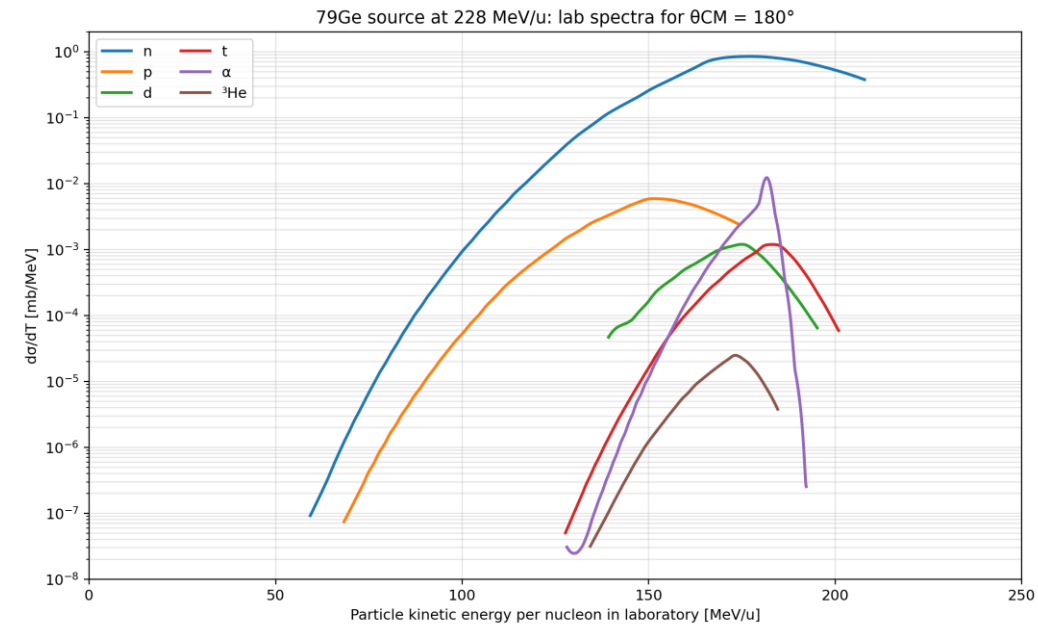
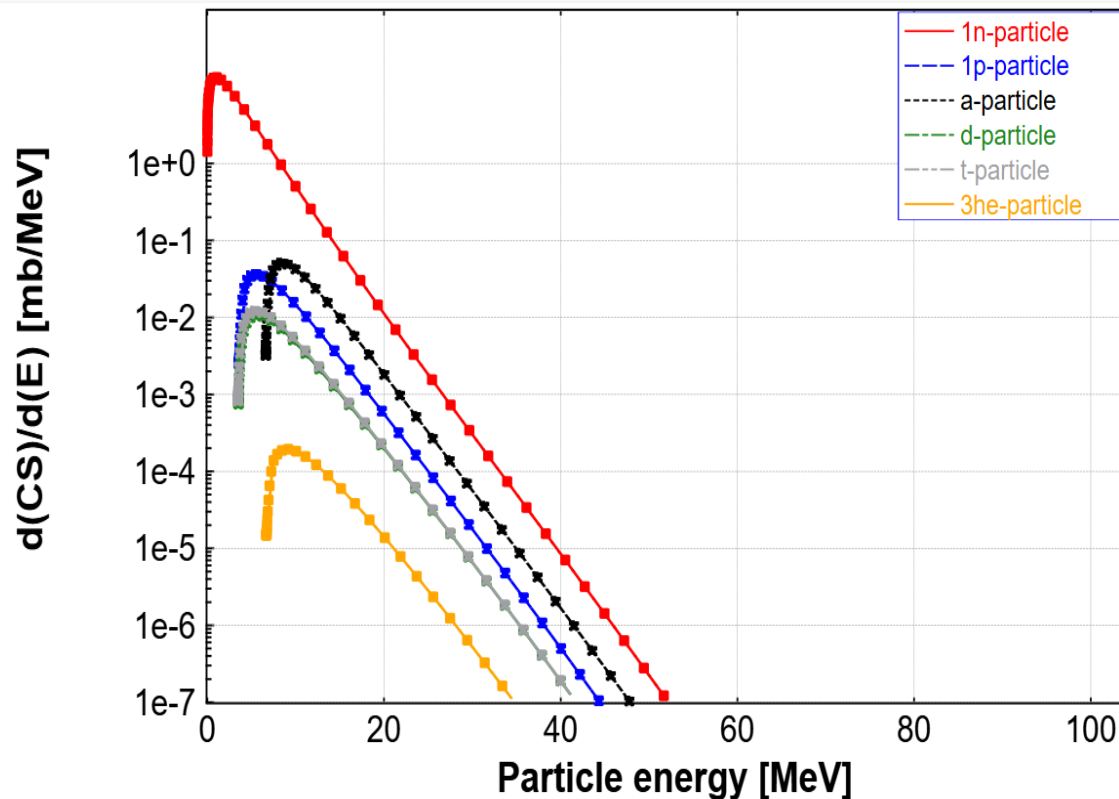
These models are developed for projectile fragmentation products, but most of light charge particles come from de-excitation

LISE++ AA

^{79}Ge excitation distributions: Emitted particle energies

ABRASION-ABLATION - $^{82}\text{Se} + \text{C}$

E* method: <3>; <T>: 16.69(0.00,0.00) MeV; No Intrin.Therm.; t_{limit} : No; DB₁=WS4_RBF; de-excitation: Q; $\delta_{\text{evap}}=7.2$; $\delta_{\text{fiss}}=14$
 NP=64; SE: "DB1+Ca13"; Den: "C"; GeomCor: "On"; Tunlg: "On"; dR=6.0fm $\text{FisBar}=\#1$; $\text{Bar}^{\text{Fac}}=1.00$; Modes=^{1010 1111 010}

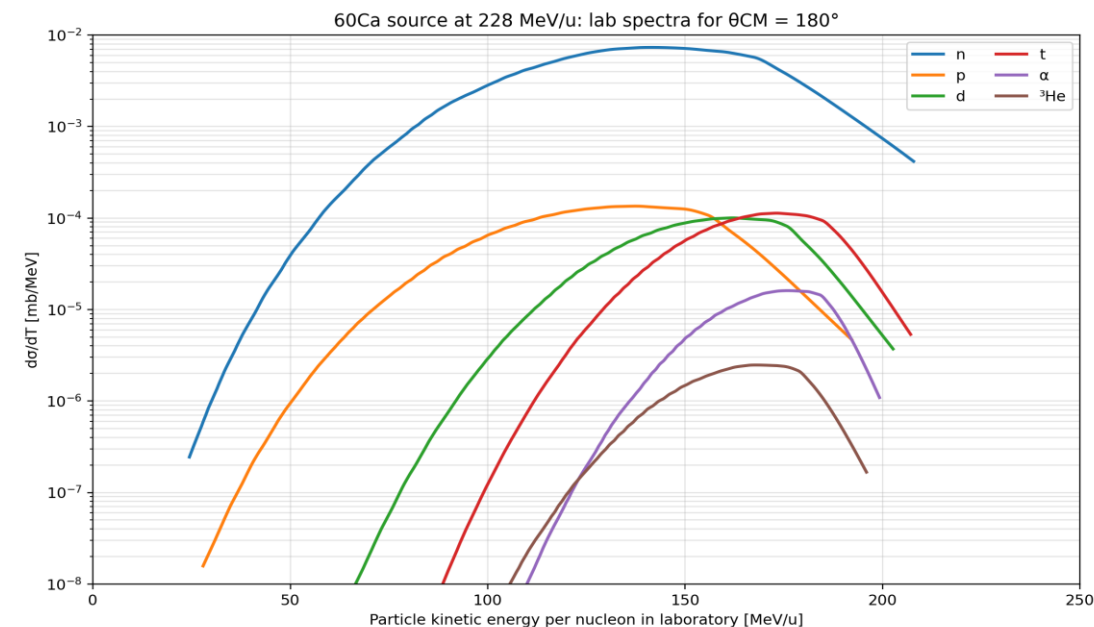
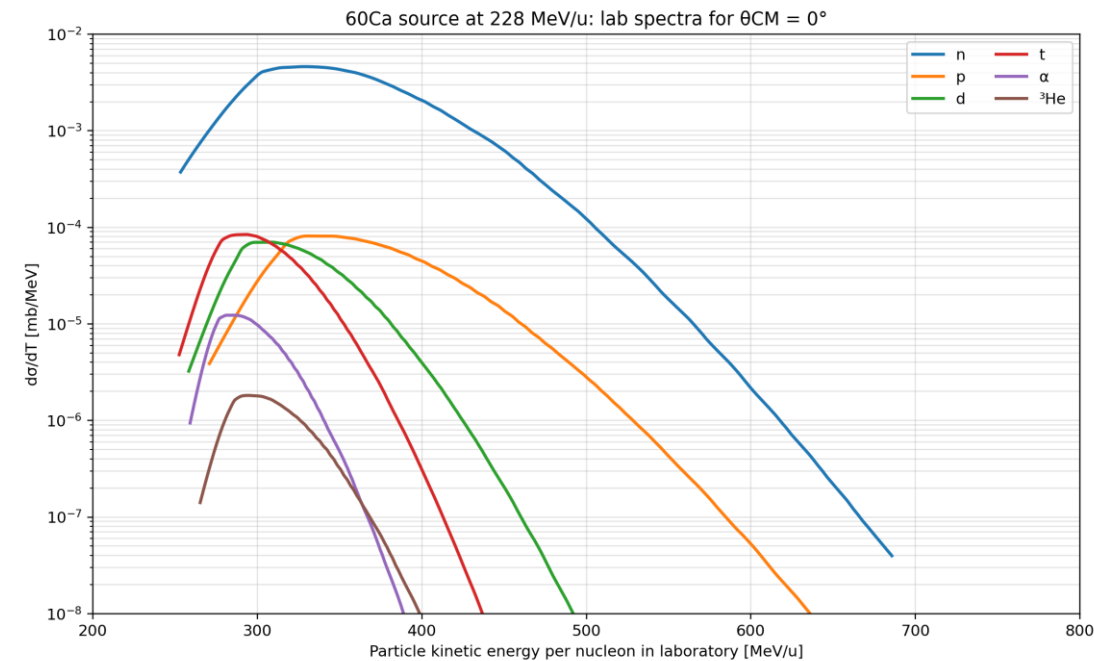
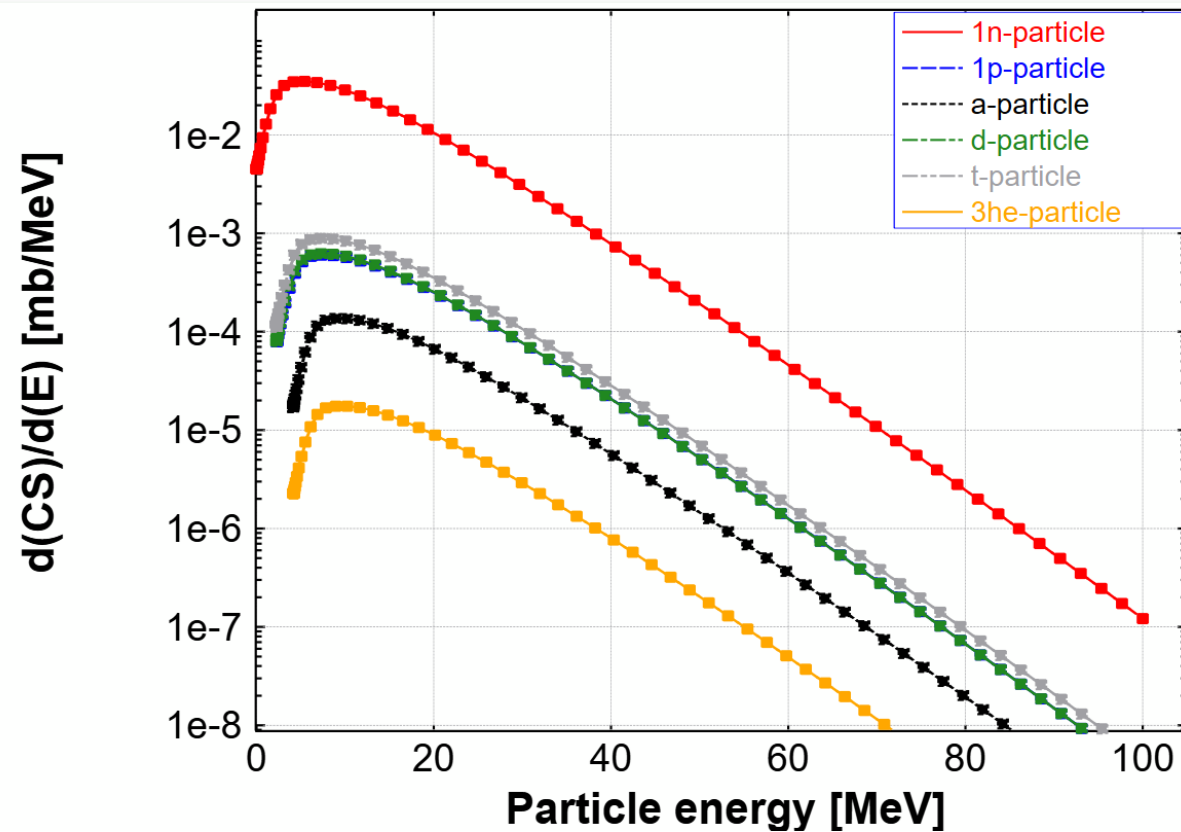


LISE⁺⁺ AA

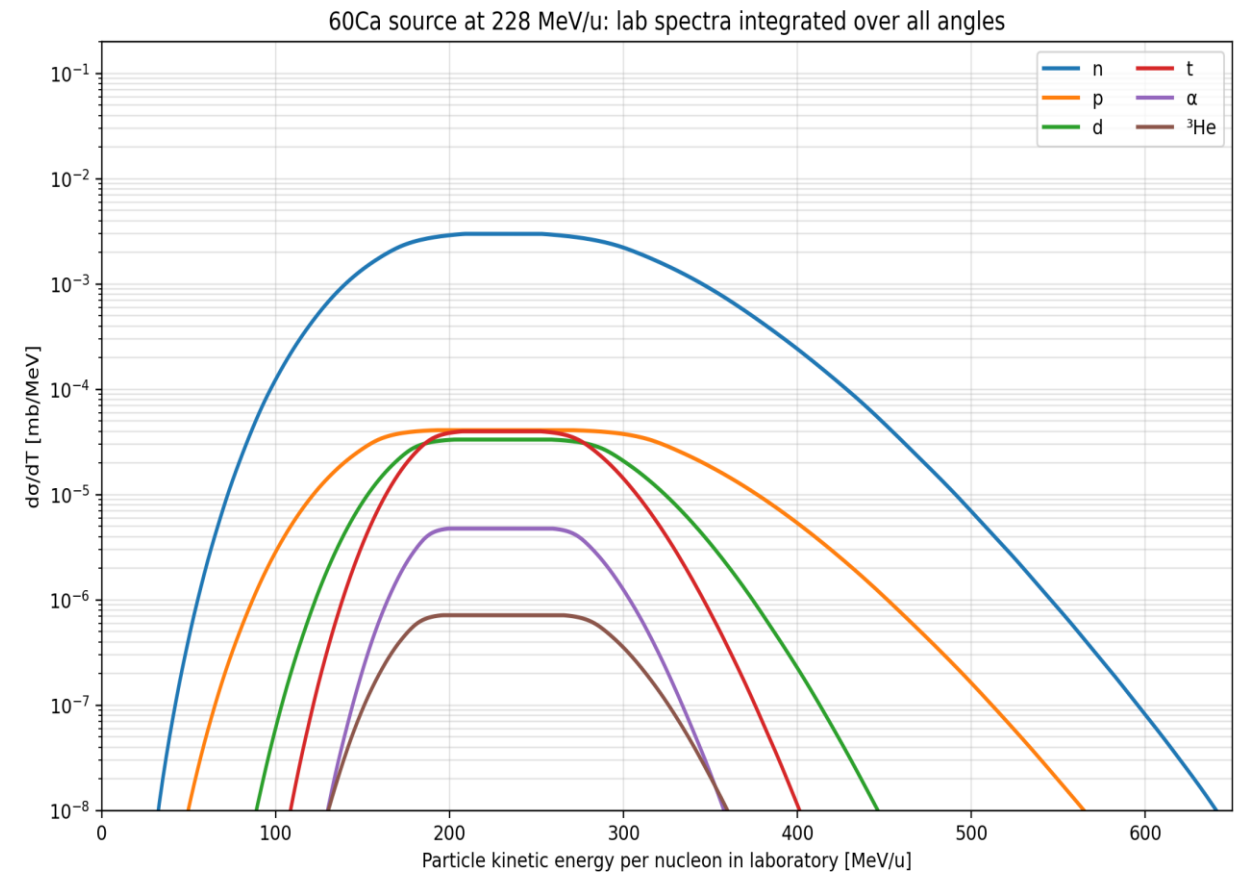
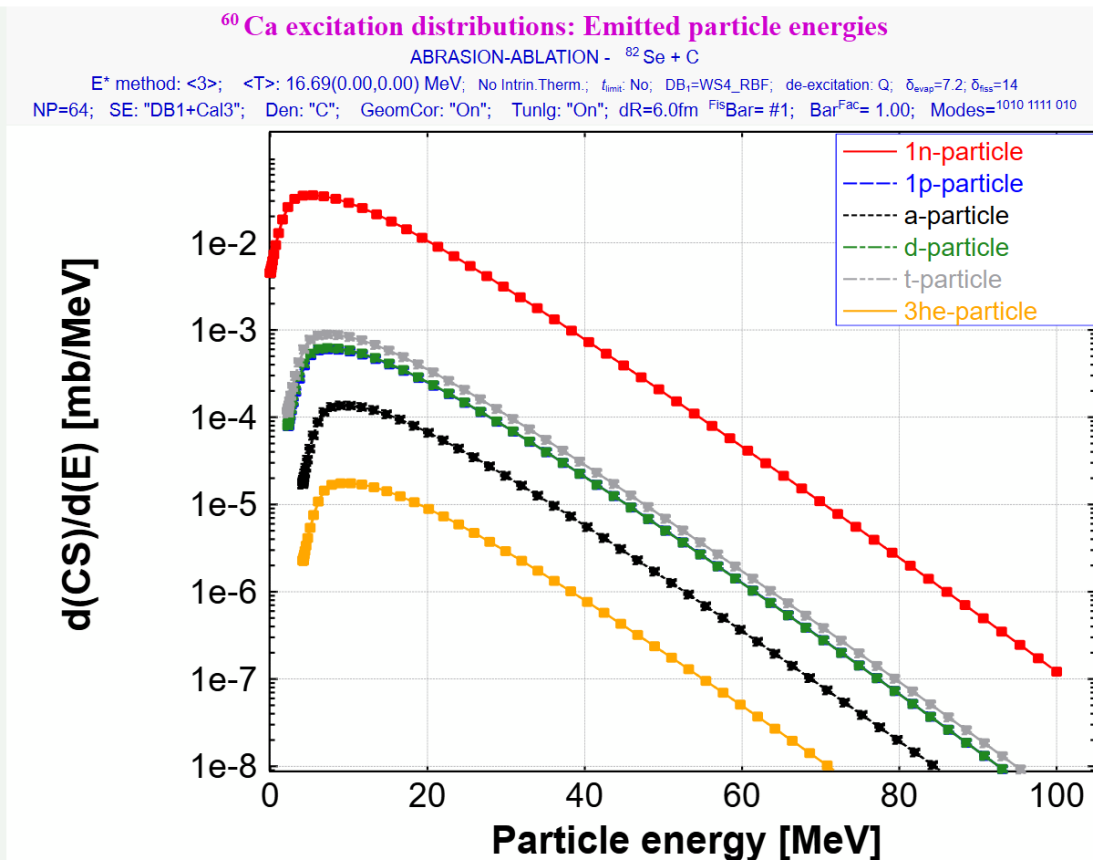
^{60}Ca excitation distributions: Emitted particle energies

ABRASION-ABLATION - $^{82}\text{Se} + \text{C}$

E* method: <3>; <T>: 16.69(0.00,0.00) MeV; No Intrin.Therm.; t_{lim} : No; DB₁=WS4_RBF; de-excitation: Q; δ_{evap} =7.2; δ_{fiss} =14
NP=64; SE: "DB1+Cal3"; Den: "C"; GeomCor: "On"; Tunlg: "On"; dR=6.0fm FisBar = #1; BarFac = 1.00; Modes=^{1010 1111 010}



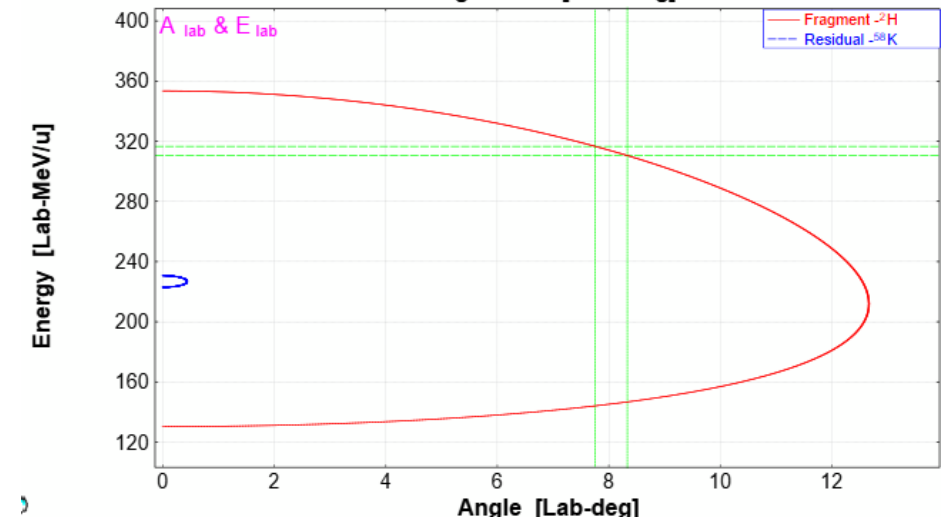
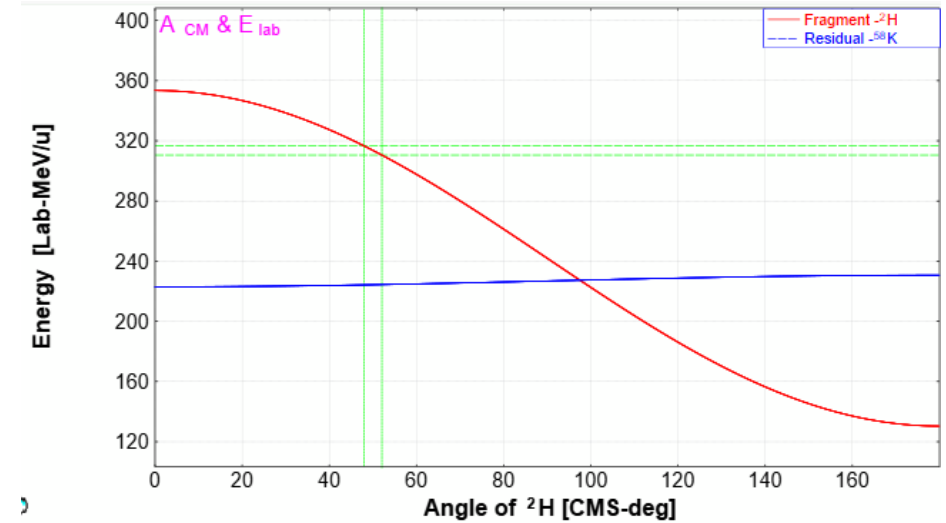
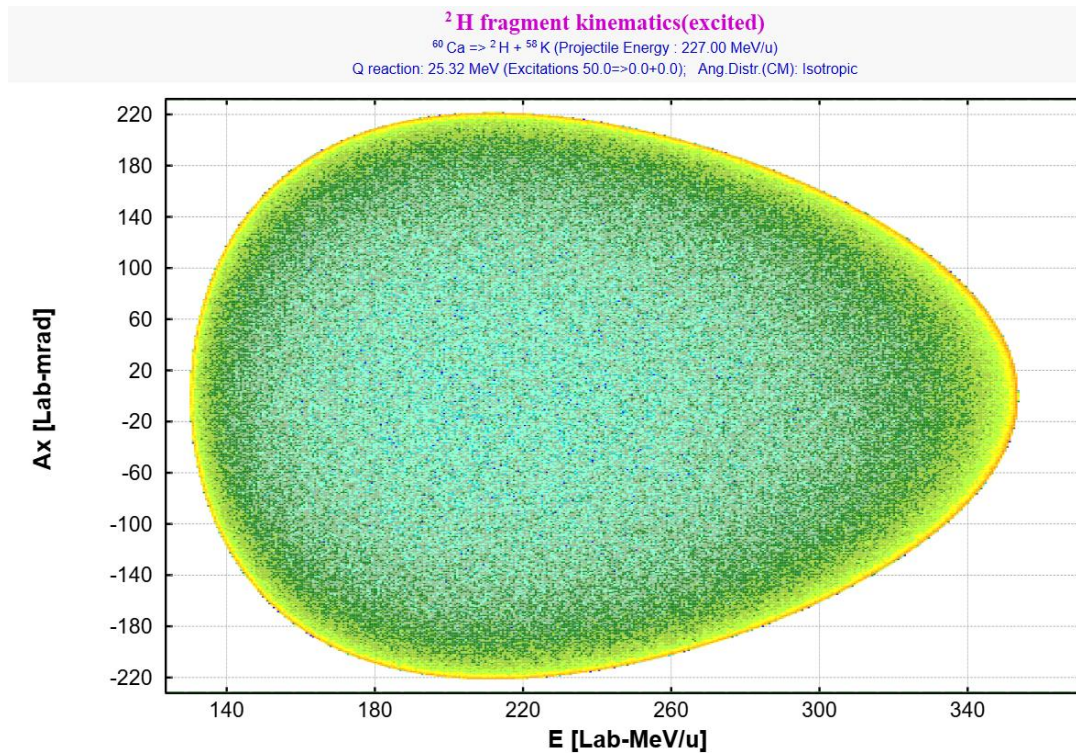
LISE⁺⁺ AA



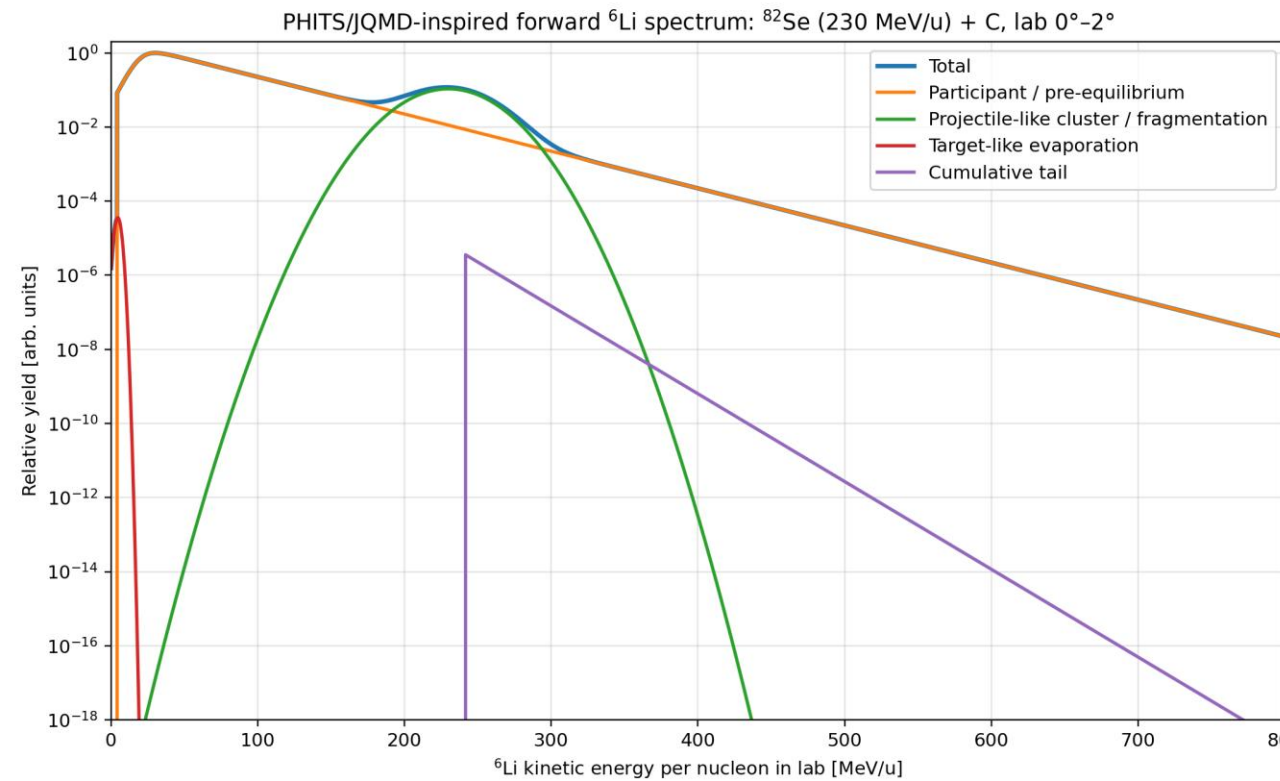
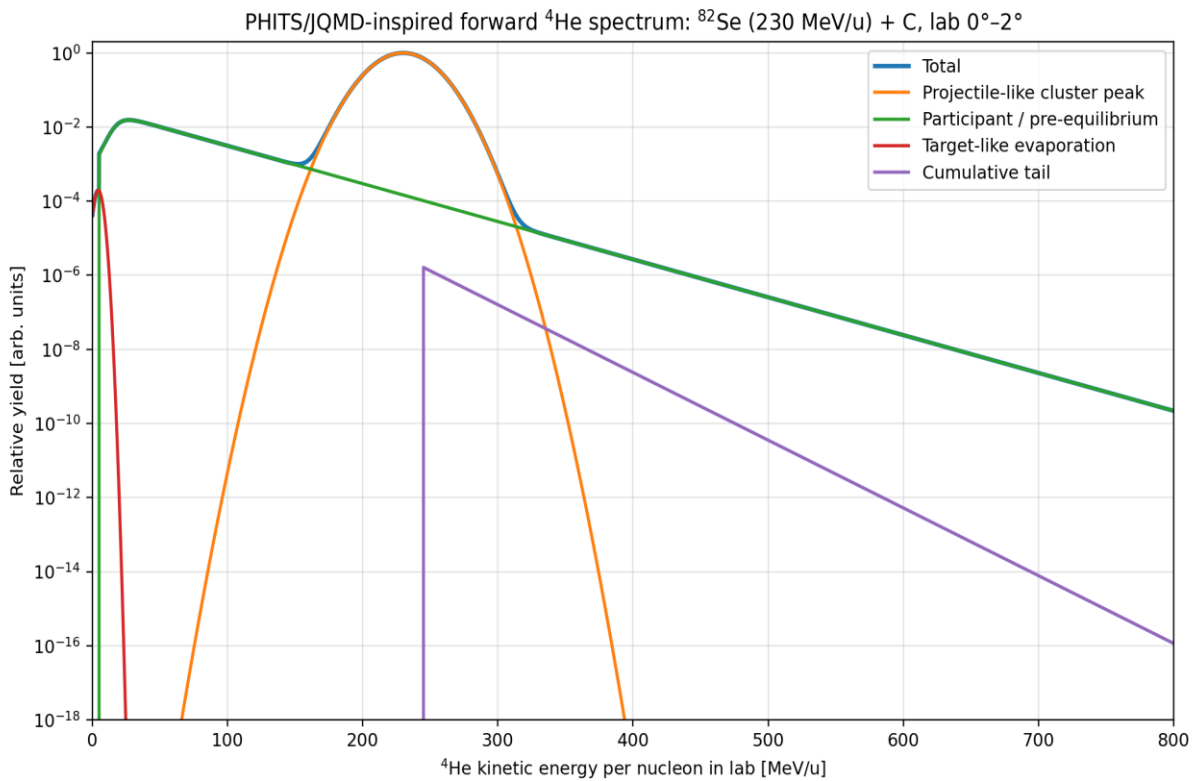
Participants

			ME [MeV]	Excitation Energy
A	Beam	⁶⁰ Ca	12.28	50
B	Target	¹² C		
C	Fragment	² H	13.14	0
D	Residual	⁵⁸ K	23.82	0.001

Deuteron emission at E*= 50 MeV

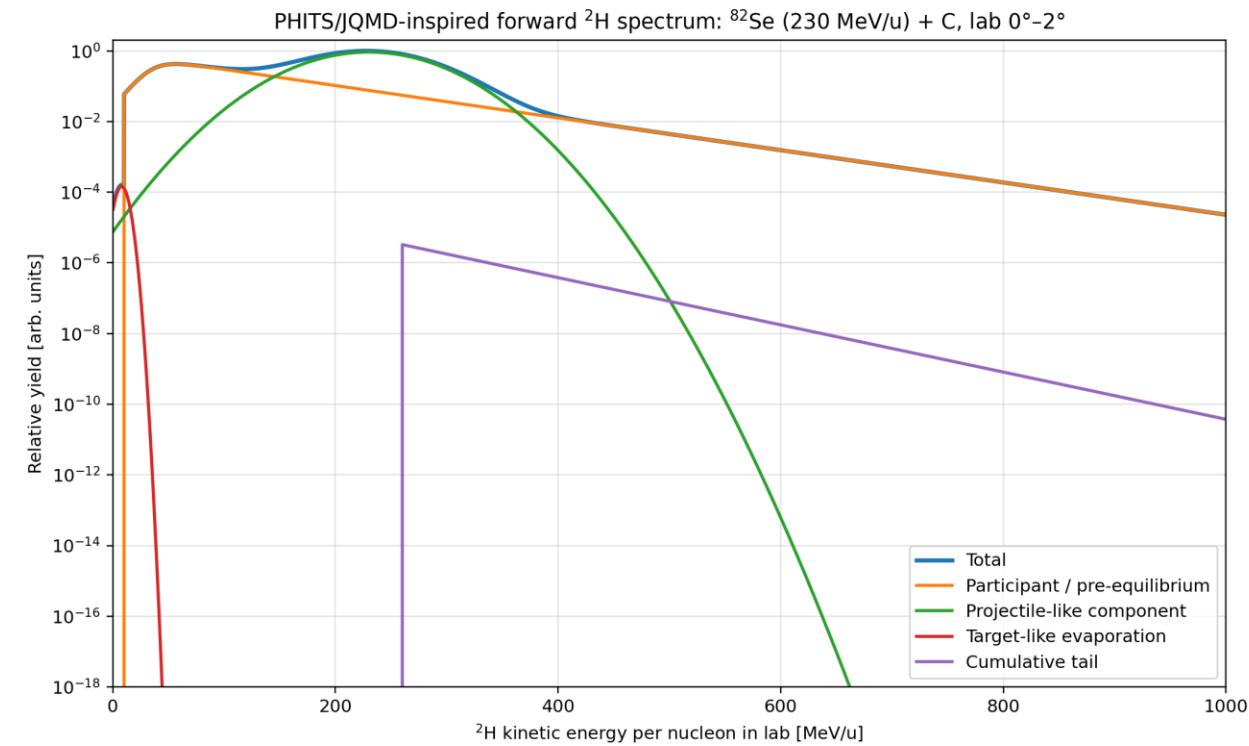
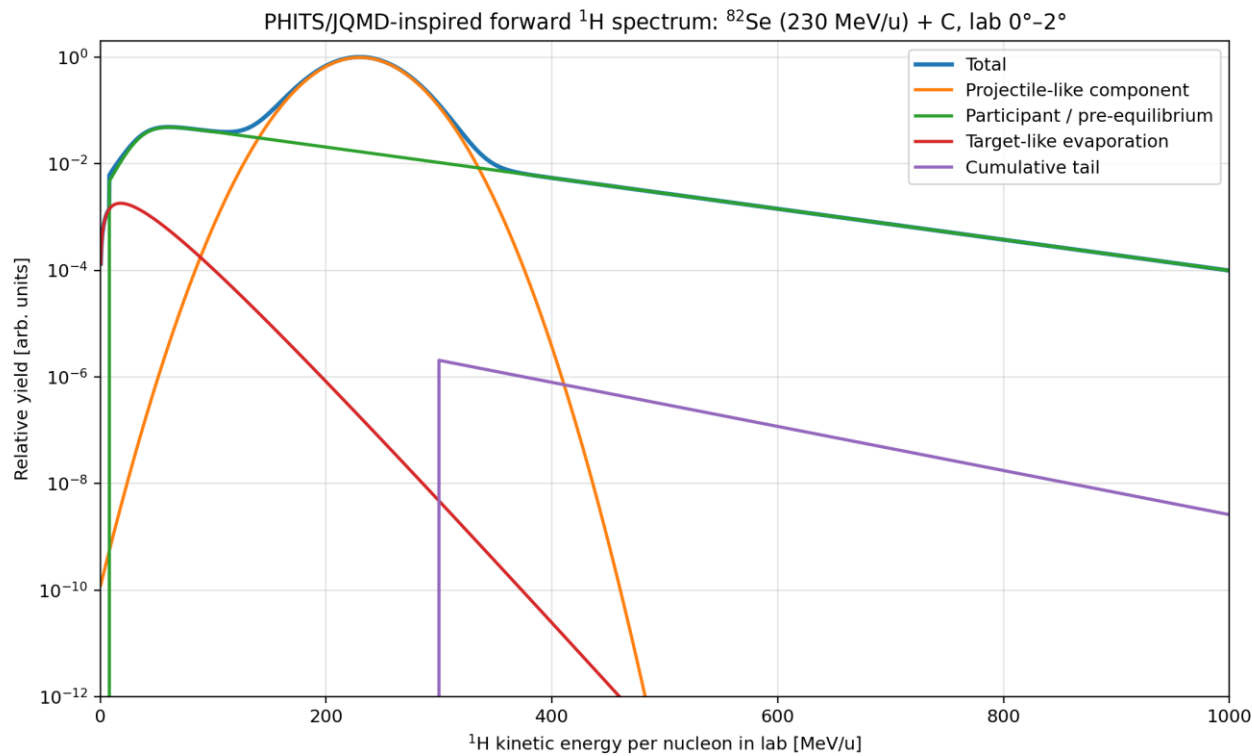


Important: Component contributions are not normalized



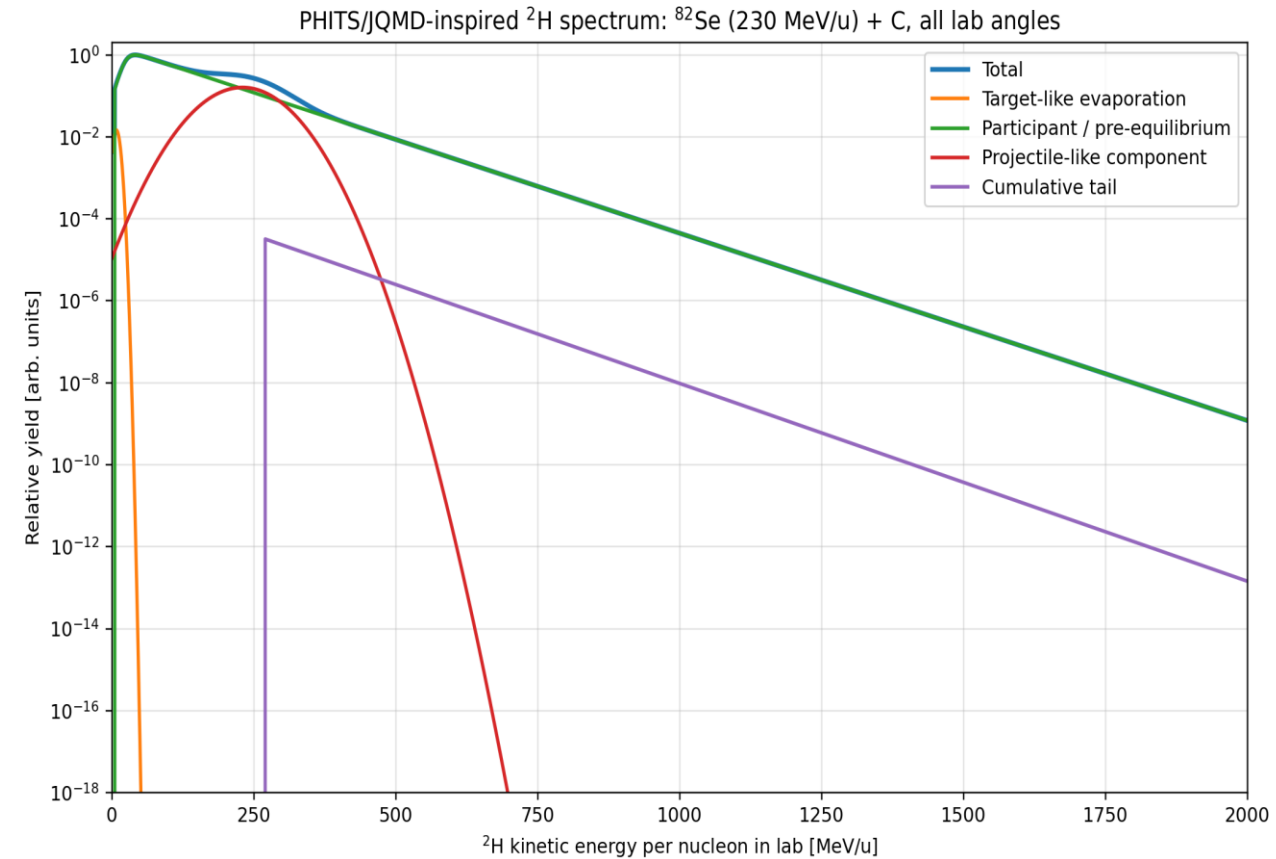
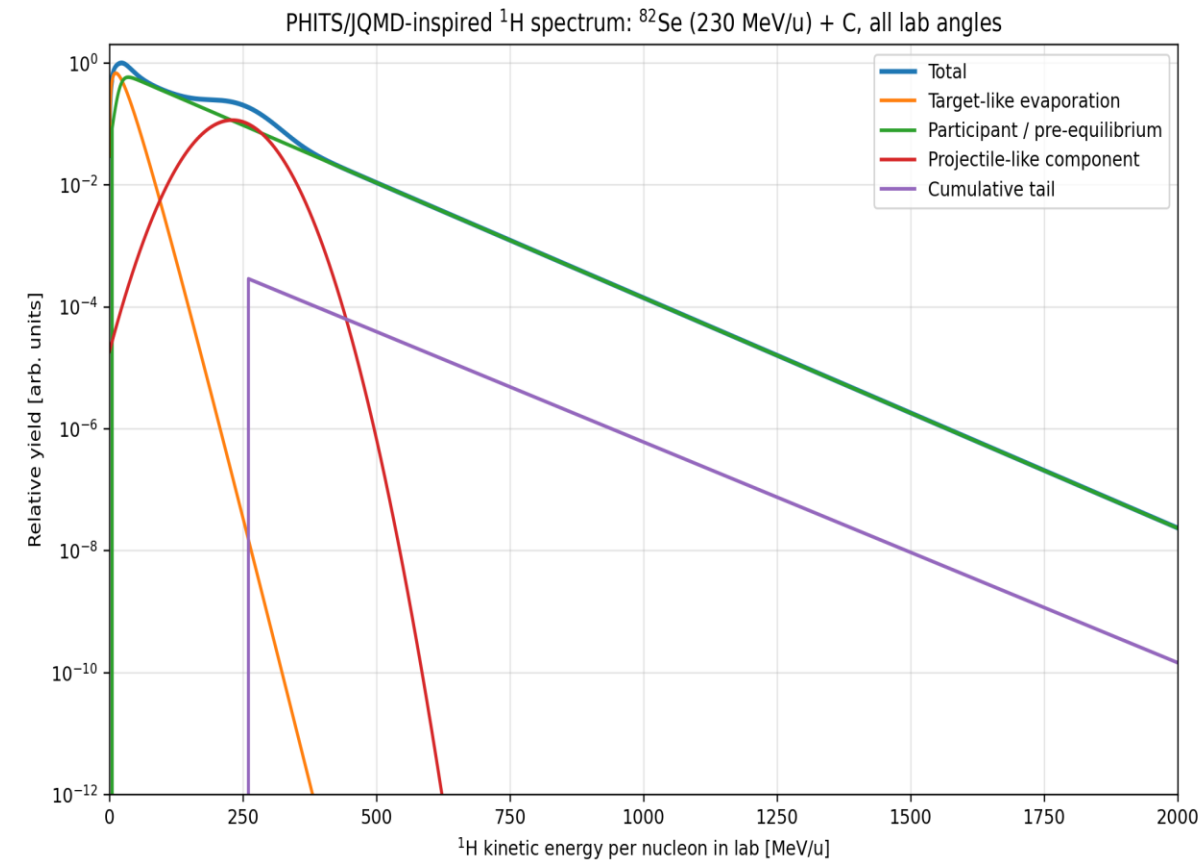
The request has been sent to Tom generate energy spectra directly from PHITS

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The request has been sent to Tom generate energy spectra directly from PHITS

Version	Topic	Implementation Summary
18.5.14	Fireball prefragment search for 1H	SearchPrefragment() now uses a temporary 2H search seed when the requested product is 1H. This makes proton momentum distributions follow the same evaporation-parent APF behavior as 2H, instead of stopping with direct-like Apf=1.

- The momentum-width investigation showed that 1H had Apf=1, while 2H, 3H, and 4He were assigned evaporation-parent-like APF values near 27-29.
- The width expression uses sigparB proportional to $\sqrt{\text{Apf} * dA / dP}$, so Apf=1 made the proton distribution significantly narrower than neighboring hydrogen isotopes.

Using a temporary 2H seed only inside SearchPrefragment() keeps the original requested product as 1H while allowing the APF search to use the same parent-search path as deuterons.

Diagnostic Values

Product	Observed Before Change	Consequence
1H	Apf=1, sigparB=194.36	Direct-like APF produced a narrow proton spectrum.
2H	Apf=27, sigparB=833.762	Evaporation-parent-like APF produced a much wider distribution.
3H	Apf=28.2208, sigparB=843.458	Width consistent with the same APF behavior as 2H.
4He	Apf=29, sigparB=848.041	Width consistent with the light-particle evaporation-parent group.

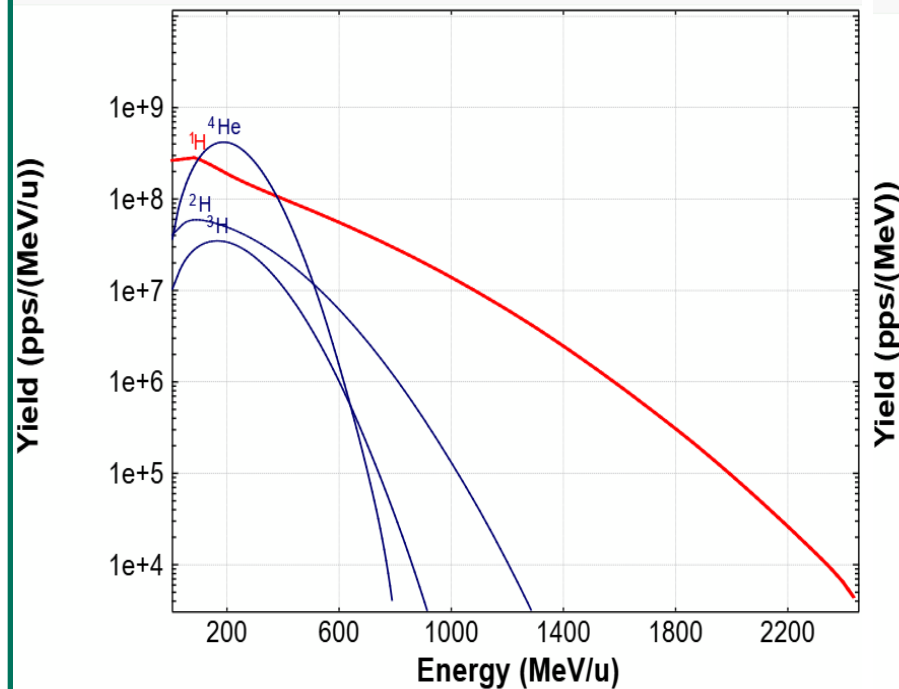
See next slide for update results

v.18.5.14, Convolution #2

Stripper → Energy: output

⁸²Se (227 MeV/u) + C (1 mm); Settings on ¹H; Config: .
dP/P=100.00%

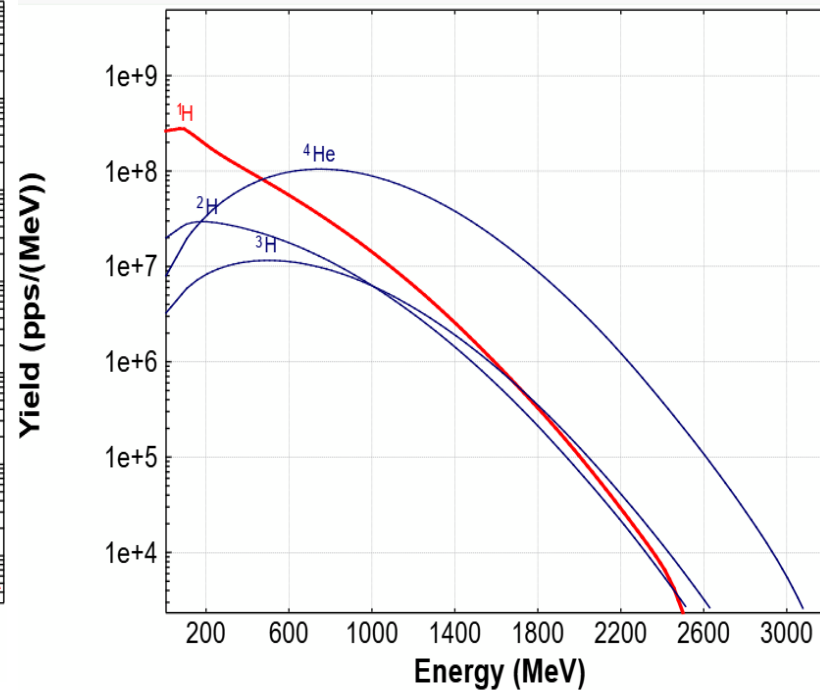
without charge states
all reactions separ.



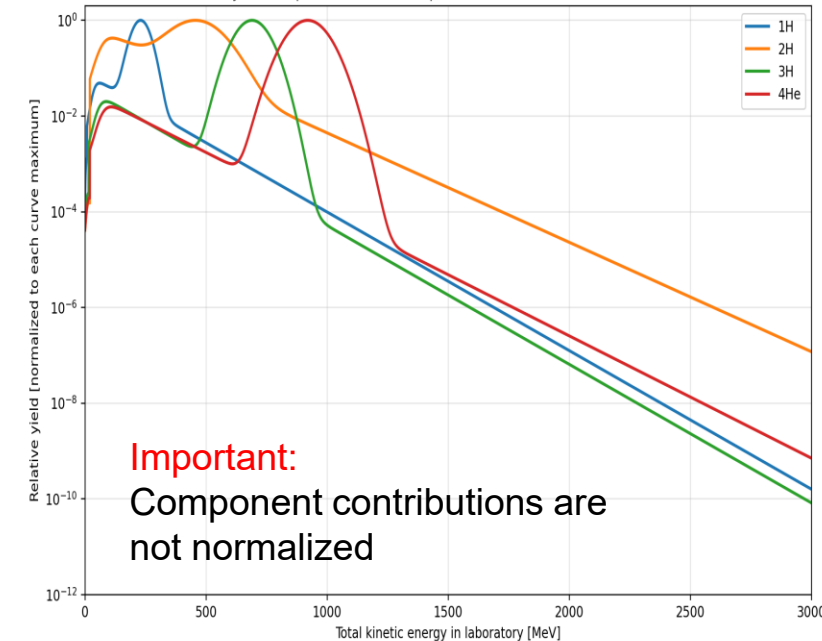
Stripper → TKE: output

⁸²Se (227 MeV/u) + C (1 mm); Settings on ¹H; Config: .
dP/P=100.00%

without charge states
all reactions separ.



PHITS/JQMD-inspired forward TKE spectra: ⁸²Se (230 MeV/u) + C, lab 0°-2°



Important:
Component contributions are
not normalized

Main recommendation

The Convolution model should not be regarded as universally superior to every other LISE⁺⁺ momentum model. However, for light particles produced predominantly through de-excitation of an excited projectile-like fragment, it is conceptually more appropriate than models developed for direct projectile-fragmentation residues.

Why the Convolution model is appropriate in this case

For an ordinary projectile-fragmentation residue, the reaction can be represented schematically as:

Projectile → excited projectile-like fragment

For a light particle emitted later from that excited fragment, there are two kinematic stages:

Projectile → F* → R + x

where x = p, d, t, ³He, α, and similar light particles.

The laboratory distribution of the emitted particle should therefore contain:

- the velocity and momentum distribution of the excited parent fragment F*;
- the particle-energy and angular distribution in the F* rest frame;
- the Lorentz transformation into the laboratory frame.

Conceptually, this is a convolution of the parent-fragment distribution with the decay kernel:

$$f_x^{\text{lab}}(p) = \int f_{F^*}(P) K_{F^* \rightarrow x}(p | P, E^*, \theta^*) dP$$

This is the central physics reason for preferring the Convolution model in the present context.

$$\begin{aligned} Y_{\text{light}} &= Y_{\text{projectile fragmentation}} \\ &+ Y_{\text{projectile de-excitation}} \\ &+ Y_{\text{participant/pre-equilibrium}} \\ &+ Y_{\text{target evaporation}} \end{aligned}$$

Recommended long-term LISE⁺⁺ solution

A more defensible treatment would separate the light-particle spectrum into mechanisms rather than assign one effective momentum model to the entire yield:

- Projectile-residue decay: parent momentum distribution convolved with the CM de-excitation spectrum;
- Direct fragmentation: Goldhaber, Morrissey, or another fragmentation model;
- Participant or pre-equilibrium emission: a dedicated broad-source distribution;
- Target evaporation: a nearly stationary and approximately isotropic source.

This mechanism-separated approach would be more physically transparent than forcing all channels into one Gaussian or one empirical longitudinal distribution.

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Important limitation

The Convolution model should not automatically replace every other model for all light particles. A calculated light-particle yield can contain several physically distinct mechanisms:

- target-like evaporation;
- projectile-residue de-excitation;
- participant or pre-equilibrium emission;
- direct breakup, fragmentation, or coalescence.

The Convolution model mainly addresses the projectile-residue de-excitation component. It does not, by itself, provide a complete description of target evaporation, fireball emission, cascade processes, or pre-equilibrium production.

Comparison with other momentum-distribution models

Goldhaber-type model

This model is well suited to the momentum width of a direct projectile-fragmentation residue, where the width is generated mainly by the random removal of projectile nucleons. It becomes questionable when applied directly to p, d, t, or α particles that are secondary products emitted after equilibration.

Morrissey-type or empirical fragmentation models

These models are useful for centroids and widths of projectile-like residues across systematic fragmentation data. They describe the transported residue, not necessarily a light particle evaporated from that residue. Direct application may therefore generate an artificial narrow structure near projectile velocity.

Phenomenological light-fragment models

These may be useful when a light isotope is produced directly through breakup, coalescence, or multifragmentation. However, no single effective model can cleanly represent direct cluster removal, participant emission, pre-equilibrium emission, statistical de-excitation, and target-like evaporation at the same time.

Convolution model

The Convolution model is the most physically motivated starting point when the dominant process is projectile fragmentation followed by de-excitation emission. It naturally broadens the laboratory distribution because the decay recoil is folded with the momentum distribution of the moving parent fragment.

Recommended model-selection statement

Use the Convolution model as the default momentum treatment for light particles identified as secondary de-excitation products of projectile-like fragments. Use direct-fragmentation models only for an explicitly direct component, and avoid interpreting one effective model as representing the complete light-particle spectrum.